

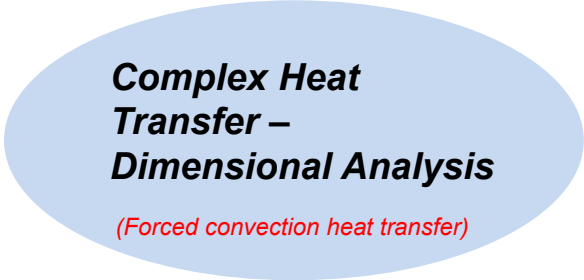


CM3110
Transport/Unit Ops I
Part II: Heat Transfer



Michigan Tech





Professor Faith Morrison

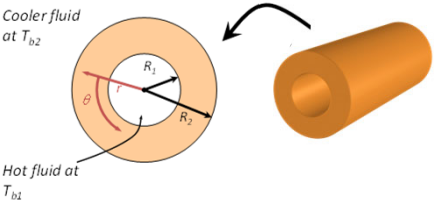
Department of Chemical Engineering
 Michigan Technological University

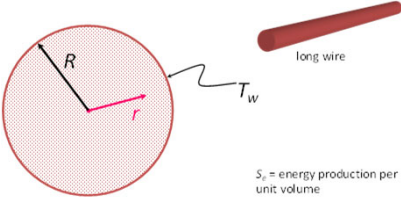
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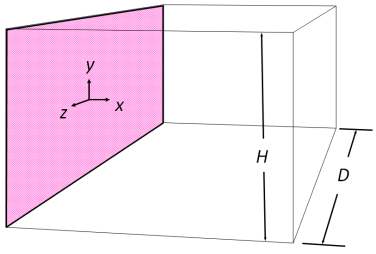
1D Heat Transfer

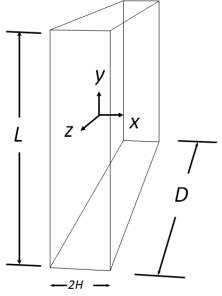
(what have we been up to?)

Examples of (simple, 1D) Heat Conduction









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1D Heat Transfer

Examples of (simple, 1D) Heat Conduction

But these are highly simplified geometries

Hot fluid at T_{b1}

Cooler fluid at T_{b2}

long wire

T_w

S_v = energy production per unit volume

H , D , $-2H$

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Complex Heat Transfer

How do we handle complex geometries, complex flows, complex machinery?

T_1 cold

T_2 less cold

T_1' less hot

T_2' hot

Q_{in}

$W_{s,on}$

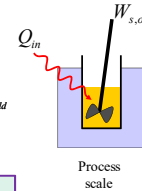
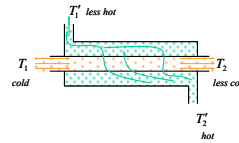
Process scale

4

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Complex Heat Transfer – Dimensional Analysis

(Answer: Use the same techniques we have been using in fluid mechanics)



Engineering Modeling

- Choose an idealized problem and solve it
- From insight obtained from **ideal** problem, identify governing equations of **real** problem
- Nondimensionalize the governing equations; deduce dimensionless scale factors (e.g. Re, Fr for fluids)
- Design experiments to test modeling thus far
- Revise modeling (structure of dimensional analysis, identity of scale factors, e.g. add roughness lengthscale)
- Design additional experiments
- Iterate until useful correlations result

5

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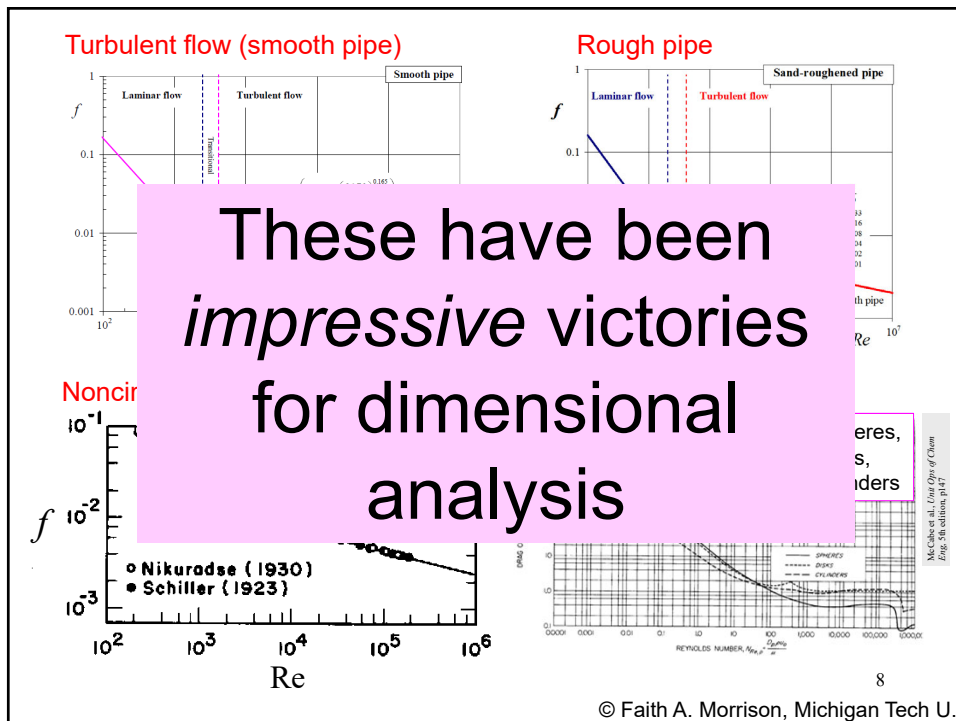
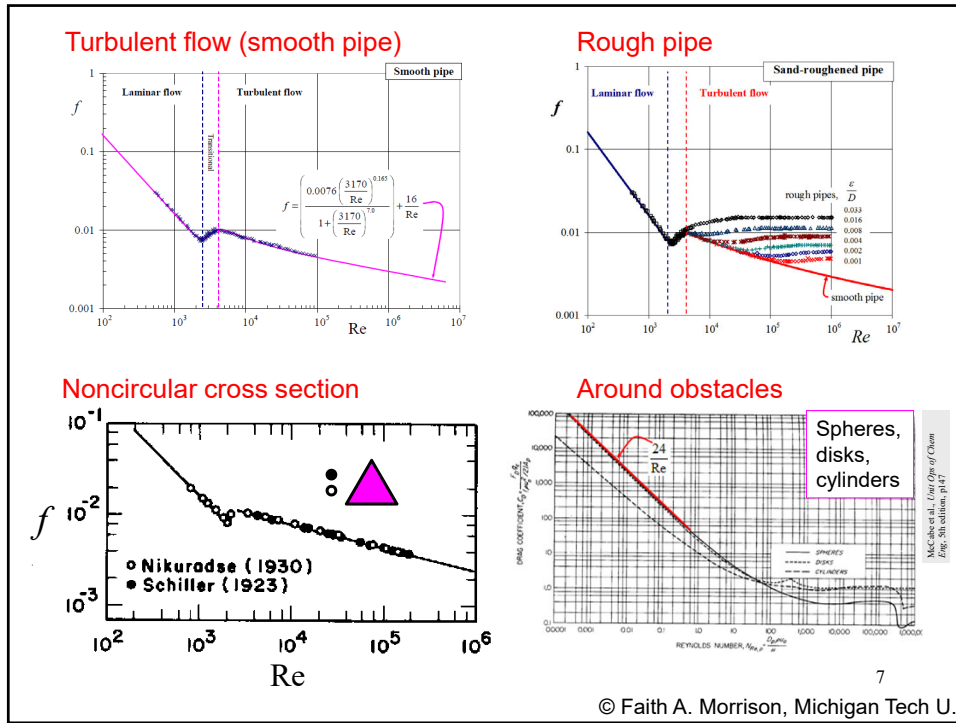
Complex Heat Transfer – Dimensional Analysis

Experience with Dimensional Analysis thus far:

- Flow in pipes at all flow rates (laminar and turbulent)
Solution: Navier-Stokes, Re, Fr, L/D , dimensionless wall force = f ; $f = f(\text{Re}, L/D)$
- Rough pipes
Solution: add additional length scale; then nondimensionalize
- Non-circular conduits
Solution: Use hydraulic diameter as the length scale of the flow to nondimensionalize
- Flow around obstacles (spheres, other complex shapes)
Solution: Navier-Stokes, Re, dimensionless drag = C_D ; $C_D = C_D(\text{Re})$
- Boundary layers
Solution: Two components of velocity need independent lengthscales

6

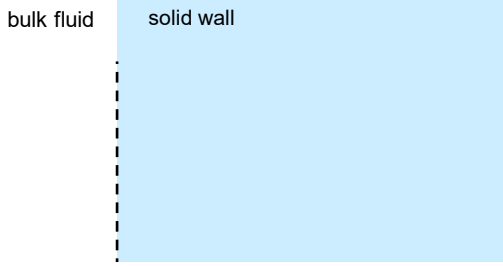
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Complex Heat Transfer – Dimensional Analysis

Now, move to heat transfer:

- Forced convection heat transfer **from fluid to wall**
Solution: ?
- Natural convection heat transfer from fluid to wall
Solution: ?
- Radiation heat transfer from solid to fluid
Solution: ?



bulk fluid solid wall

9

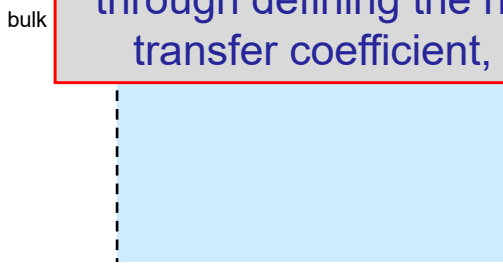
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Complex Heat Transfer – Dimensional Analysis

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- Forced convection heat transfer from fluid to wall
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Solution: ?
- Radiation heat transfer from solid to fluid
Solution: ?

We have already started using the results/techniques of dimensional analysis through defining the heat transfer coefficient, h



bulk solid wall

10

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Complex Heat Transfer – Dimensional Analysis

Now, move to heat transfer:

- Forced convection heat transfer from fluid to wall
Solution: ?
- Natural convection heat transfer from fluid to wall
Solution: ?
- Radiation heat transfer from fluid to wall
Solution: ?

bulk

We have already started using the results/techniques of dimensional analysis through defining the heat transfer coefficient, h

(recall that we did this in fluids too: we used the $f(Re)$ correlation (MEB+Moody chart) long before we knew where that all came from)

11

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Handy tool: Heat Transfer Coefficient

The temperature variation in fluid near-wall region is caused by complex phenomena that are lumped together into the heat transfer coefficient, h

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Complex Heat Transfer – Dimensional Analysis

The flux at the wall is given by the empirical expression known as **Newton's Law of Cooling**

This expression serves as the definition of the **heat transfer coefficient**.

$$\left| \frac{q_x}{A} \right| \equiv h |T_{bulk} - T_{wall}|$$

***h* depends on:**

- geometry
- fluid velocity
- fluid properties
- temperature difference

13

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Complex Heat Transfer – Dimensional Analysis

The flux at the wall is given by the empirical expression known as **Newton's Law of Cooling**

This expression serves as the definition of the **heat transfer coefficient**.

$$\left| \frac{q_x}{A} \right| \equiv h |T_{bulk} - T_{wall}|$$

To get values of ***h*** for various situations, we need to measure data and create data correlations (**dimensional analysis**)

***h* depends on:**

- geometry
- fluid velocity
- fluid properties
- temperature difference

14

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Complex Heat Transfer – Dimensional Analysis

Complex Heat transfer Problems to Solve:

- Forced convection heat transfer from fluid to wall
Solution: ?
- Natural convection heat transfer from fluid to wall
Solution: ?
- Radiation heat transfer from solid to fluid
Solution: ?

- The functional form of h will be different for these **three** situations (different physics)
- Investigate simple problems in each category, model them, take data, correlate

15
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Complex Heat Transfer – Dimensional Analysis

Complex Heat transfer Problems to Solve:

- Forced convection heat transfer from fluid to wall
Solution: ?
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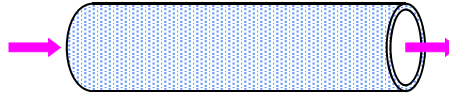
Let's look at forced convection in a pipe. There are three pieces to the **physics**:

- Pipe flow
- Energy
- Boundary conditions

16
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Complex Heat Transfer – Dimensional Analysis

Chosen problem: Forced Convection Heat Transfer
Solution: Dimensional Analysis



Following procedure familiar from pipe flow,

- **What are governing equations?**
- **Scale factors (dimensionless numbers)?**
- **Quantity of interest?**

Answer: Heat flux at the wall

17

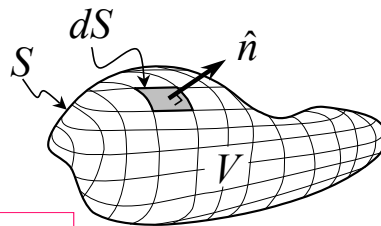
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Complex Heat Transfer – Dimensional Analysis

General Energy Transport Equation

(microscopic energy balance)

As for the derivation of the microscopic momentum balance, the microscopic energy balance is derived on an arbitrary volume, V , enclosed by a surface, S .



Gibbs notation:

$$\rho \hat{C}_p \left(\frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S$$

see handout for
component notation

18

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Complex Heat Transfer – Dimensional Analysis

General Energy Transport Equation

(microscopic energy balance; **in the fluid**)

$$\rho \hat{C}_p \left(\underbrace{\frac{\partial T}{\partial t}}_{\text{rate of change}} + \underbrace{\underline{v} \cdot \nabla T}_{\text{convection}} \right) = \underbrace{k \nabla^2 T}_{\text{conduction (all directions)}} + \underbrace{S}_{\text{source (energy generated per unit volume per time)}}$$

fluid velocity must satisfy equation of motion, equation of continuity

see handout for component notation

19

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Complex Heat Transfer – Dimensional Analysis

The Equation of Energy for systems with constant k

Microscopic energy balance, constant thermal conductivity; Gibbs notation

$$\rho \hat{C}_p \left(\frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S$$

Microscopic energy balance, constant thermal conductivity; Cartesian coordinates

$$\rho \hat{C}_p \left(\frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = k \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

Microscopic energy balance, constant thermal conductivity; cylindrical coordinates

$$\rho \hat{C}_p \left(\frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

Microscopic energy balance, constant thermal conductivity; spherical coordinates

$$\rho \hat{C}_p \left(\frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial T}{\partial \phi} \right) = k \left(\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial T}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial T}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 T}{\partial \phi^2} \right) + S$$

Note: this handout is also on the web

<https://pages.mtu.edu/~fmorriso/cm310/energy.pdf>

20

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**** REVIEW ** REVIEW ****

Example: Heat flux in a cylindrical shell

Assumptions:

- long pipe
- steady state
- k = thermal conductivity of wall
- h_1, h_2 = heat transfer coefficients

*What is the steady state temperature profile in a cylindrical shell (pipe) if the **fluid on the inside** is at T_{b1} and the **fluid on the outside** is at T_{b2} ? ($T_{b1} > T_{b2}$)*

Forced-convection heat transfer

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Now: How do develop correlations for h ?

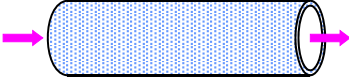
Consider: Heat-transfer to/from flowing fluid inside of a tube – forced-convection heat transfer

T_1 = core bulk temperature
 T_o = wall temperature
 $T(r, \theta, z)$ = temp distribution in the fluid

Instead:
Dimensional
Analysis and Data
correlations

In principle, with the right math/computer tools, we could calculate the complete temperature and velocity profiles in the moving fluid.

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Dimensional Analysis in Forced Convection Heat Transfer 

Pipe flow

z-component of the Navier-Stokes Equation:

$$\rho \left(\frac{\partial v_z}{\partial t} + v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z} \right) = -\frac{\partial P}{\partial z} + \mu \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 v_z}{\partial \theta^2} + \frac{\partial^2 v_z}{\partial z^2} \right) + \rho g_z$$

Choose:

D = characteristic length

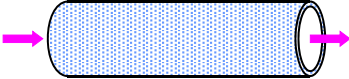
V = characteristic velocity

D/V = characteristic time

ρV^2 = characteristic pressure

- Choose “typical” values (*scale factors*)
- Use them to scale the equations
- Deduce which terms dominate

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Dimensional Analysis in Forced Convection Heat Transfer 

Pipe flow

non-dimensional variables:

time:

$$t^* \equiv \frac{tV}{D}$$

position:

$$r^* \equiv \frac{r}{D}$$

$$z^* \equiv \frac{z}{D}$$

velocity:

$$v_z^* \equiv \frac{v_z}{V}$$

$$v_r^* \equiv \frac{v_r}{V}$$

$$v_\theta^* \equiv \frac{v_\theta}{V}$$

driving force:

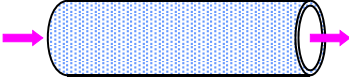
$$P^* \equiv \frac{P}{\rho V^2}$$

$$g_z^* \equiv \frac{g_z}{g}$$

- Choose “typical” values (*scale factors*)
- Use them to scale the equations
- Deduce which terms dominate

24
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Dimensional Analysis in Forced Convection Heat Transfer



Energy

Microscopic energy balance:

$$\rho \hat{c}_p \left(\frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

non-dimensional variables:

<p>position:</p> $r^* \equiv \frac{r}{D}$ $z^* \equiv \frac{z}{D}$	<p>temperature:</p> $T^* \equiv \frac{T - T_o}{(T_1 - T_o)}$	<p>source:</p> $S^* \equiv \frac{S}{S_0}$
---	---	--


Choose:
 T – use a characteristic **interval** (since distance from absolute zero is not part of this physics)
 S – use a reference source, S_0

$$S_0 \equiv \frac{(T_1 - T_o) V \rho \hat{c}_p}{D} \left[\frac{W}{m^3} \right]$$

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Complex Heat Transfer – Dimensional Analysis Forced Convection

Forced Convection Heat Transfer



Pipe flow

non-dimensional variables:

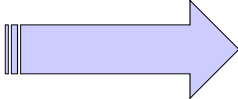
<p>time:</p> $t^* \equiv \frac{tV}{D}$	<p>position:</p> $r^* \equiv \frac{r}{D}$ $z^* \equiv \frac{z}{D}$	<p>velocity:</p> $v_z^* \equiv \frac{v_z}{V}$ $v_r^* \equiv \frac{v_r}{V}$ $v_\theta^* \equiv \frac{v_\theta}{V}$	<p>driving force:</p> $P^* \equiv \frac{P}{\rho V^2}$ $g_z^* \equiv \frac{g_z}{g}$
---	---	--	---

- Choose "typical" values (scale factors)
- Use them to scale the equations
- Deduce which terms dominate


Substitute all these definitions,

$$(t^*, r^*, z^*, p^*, g^*, v_r^*, v_\theta^*, v_z^*, T^*, S^*)$$

into the **governing equations** and simplify...



Forced Convection Heat Transfer



Energy

Microscopic energy balance:

$$\rho \hat{c}_p \left(\frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

non-dimensional variables:

<p>position:</p> $r^* \equiv \frac{r}{D}$ $z^* \equiv \frac{z}{D}$	<p>temperature:</p> $T^* \equiv \frac{T - T_o}{(T_1 - T_o)}$	<p>source:</p> $S^* \equiv \frac{S}{S_0}$
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Choose:
 T – use a characteristic **interval** (since distance from absolute zero is not part of this physics)
 S – use a reference source, S_0

$$S_0 \equiv \frac{(T_1 - T_o) V \rho \hat{c}_p}{D} \left[\frac{W}{m^3} \right]$$

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Complex Heat Transfer – Dimensional Analysis Forced Convection

Non-dimensional Energy Equation

$$\left(\frac{\partial T^*}{\partial t^*} + v_r^* \frac{\partial T^*}{\partial r^*} + \frac{v_\theta^*}{r^*} \frac{\partial T^*}{\partial \theta} + v_z^* \frac{\partial T^*}{\partial z^*}\right) = \frac{1}{\text{Pe}} \left(\frac{1}{r^*} \frac{\partial}{\partial r^*} \left(r^* \frac{\partial T^*}{\partial r^*} \right) + \frac{1}{r^{*2}} \frac{\partial^2 T^*}{\partial \theta^2} + \frac{\partial^2 T^*}{\partial z^{*2}} \right) + S^*$$

Non-dimensional Navier-Stokes Equation

$$\frac{Dv_z^*}{Dt^*} = -\frac{\partial P^*}{\partial z^*} + \frac{1}{\text{Re}} (\nabla^2 v_z^*) + \frac{1}{\text{Fr}} g^*$$

$\text{Pe} = \text{PrRe} = \frac{\hat{C}_p \mu}{k} \frac{\rho V D}{\mu}$
 $\text{Pr} = \frac{\hat{C}_p \mu}{k}$

Non-dimensional Continuity Equation

$$\frac{\partial v_x^*}{\partial x^*} + \frac{\partial v_y^*}{\partial y^*} + \frac{\partial v_z^*}{\partial z^*} = 0$$

$T^* = T^*(\text{Re}, \text{Pr})$
 $\underline{v}^* = \underline{v}^*(\text{Re}, \text{Fr})$

$\frac{Dv_z}{Dt} \equiv \frac{\partial v_z}{\partial t} + v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z}$

27
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Complex Heat Transfer – Dimensional Analysis Forced Convection

How do we get to the “*engineering quantities of interest*” ?

What are the “*engineering quantities of interest*” ?

Complex Heat Transfer – Dimensional Analysis

Non-dimensional Energy Equation

$$\left(\frac{\partial T^*}{\partial t^*} + v_r^* \frac{\partial T^*}{\partial r^*} + \frac{v_\theta^*}{r^*} \frac{\partial T^*}{\partial \theta} + v_z^* \frac{\partial T^*}{\partial z^*}\right) = \frac{1}{\text{Pe}} \left(\frac{1}{r^*} \frac{\partial}{\partial r^*} \left(r^* \frac{\partial T^*}{\partial r^*} \right) + \frac{1}{r^{*2}} \frac{\partial^2 T^*}{\partial \theta^2} + \frac{\partial^2 T^*}{\partial z^{*2}} \right) + S^*$$

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Non-dimensional Continuity Equation

$$\frac{\partial v_x^*}{\partial x^*} + \frac{\partial v_y^*}{\partial y^*} + \frac{\partial v_z^*}{\partial z^*} = 0$$

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Microscopic Energy Balance Review—Boundary Conditions

The interface between the solid and the fluid calls for a new type of boundary condition, **Newton's Law of Cooling**.

$Q = q_x$ is the engineering quantity of interest

$$\left| \frac{q_x}{A} \right| = h |T_{bulk} - T_{wall}|$$

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Forced Convection Heat Transfer

Linear driving force model $\left| \frac{q_x}{A} \right| = h |T_1 - T_0|$

Apply at the surface (in fluid):

$$(2\pi RL)(h)(T_1 - T_0) = Q = \iint_S [\hat{e}_r \cdot \tilde{q}]_{surface} dS$$

$$(2\pi RL)(h)(T_1 - T_0) = Q = \int_0^{2\pi} \int_0^L -k \frac{\partial T}{\partial r} \Big|_{r=R} R dz d\theta$$

Now, non-dimensionalize this expression as well.

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Complex Heat Transfer – Dimensional Analysis

Non-dimensionalize

non-dimensional variables:

position:

$$r^* \equiv \frac{r}{D}$$

$$z^* \equiv \frac{z}{D}$$

temperature:

$$T^* \equiv \frac{T - T_o}{(T_1 - T_o)}$$

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Complex Heat Transfer – Dimensional Analysis

$$h(\cancel{\pi DL})(\cancel{T_1 - T_o}) = \int_0^{2\pi} \int_0^{L/D} -k \frac{\partial T^*}{\partial r^*} \Big|_{r^*=1/2} \frac{(\cancel{T_1 - T_o}) \cancel{D^2}}{\cancel{D}} dz^* d\theta$$

$$2\pi \left(\frac{hD}{k} \right) \left(\frac{L}{D} \right) = \int_0^{2\pi} \int_0^{L/D} - \frac{\partial T^*}{\partial r^*} \Big|_{r^*=1/2} dz^* d\theta$$

Nusselt number, Nu
(dimensionless heat-transfer coefficient)

$$Nu = Nu\left(T^*, \frac{L}{D}\right)$$

one additional dimensionless group

32
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Complex Heat Transfer – Dimensional Analysis

$$h(\pi DL)(T_1 - T_o) = \int_0^{2\pi} \int_0^{L/D} -k \frac{\partial T^*}{\partial r^*} \Big|_{r^*=1/2} \frac{(T_1 - T_o) D^2}{D} \frac{D^2}{2} dz^* d\theta$$

This is a function of Re and Pr through fluid v_z^* distribution and energy balance

$$2\pi \left(\frac{hD}{k} \right) \left(\frac{L}{D} \right) = \int_0^{2\pi} \int_0^{L/D} \frac{\partial T^*}{\partial r^*} \Big|_{r^*=1/2} dz^* d\theta$$

Nusselt number, Nu
(dimensionless heat-transfer coefficient)

$$Nu = Nu \left(T^*, \frac{L}{D} \right)$$

one additional dimensionless group

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Complex Heat Transfer – Dimensional Analysis

Non-dimensional Energy Equation

$$\left(\frac{\partial T^*}{\partial t^*} + v_r^* \frac{\partial T^*}{\partial r^*} + \frac{v_\theta^*}{r^*} \frac{\partial T^*}{\partial \theta} + v_z^* \frac{\partial T^*}{\partial z^*} \right) = \frac{1}{Pe} \left(\frac{1}{r^*} \frac{\partial}{\partial r^*} \left(r^* \frac{\partial T^*}{\partial r^*} \right) + \frac{1}{r^{*2}} \frac{\partial^2 T^*}{\partial \theta^2} + \frac{\partial^2 T^*}{\partial z^{*2}} \right) + S^*$$

Non-dimensional Navier-Stokes Equation

$$\frac{Dv_z^*}{Dt} = -\frac{\partial P^*}{\partial z^*} + \frac{1}{Re} (\nabla^2 v_z^*) + \frac{1}{Fr} g^*$$

Non-dimensional Continuity Equation

$$\frac{\partial v_x^*}{\partial x^*} + \frac{\partial v_y^*}{\partial y^*} + \frac{\partial v_z^*}{\partial z^*} + = 0$$

Quantity of interest

$$Nu = \frac{1}{2\pi L / D} \int_0^{2\pi} \int_0^{L/D} \frac{\partial T^*}{\partial r^*} \Big|_{r^*=1/2} dz^* d\theta$$

$Pe = Pr Re = \frac{\hat{C}_p \mu}{k} \frac{\rho V D}{\mu}$
 $Pr = \frac{\hat{C}_p \mu}{k}$

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Complex Heat Transfer – Dimensional Analysis

According to our **dimensional analysis** calculations, the dimensionless heat transfer coefficient should be found to be a function of ~~four~~ dimensionless groups:

~~three~~

no free surfaces

Peclet number

$$Pe \equiv \frac{\rho \hat{c}_p V D}{k} = \frac{\hat{c}_p \mu \rho V D}{k \mu}$$

$$Nu = Nu \left(Re, Pr, \cancel{Fr}, \frac{L}{D} \right)$$

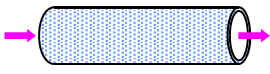
Prandtl number

$$Pr \equiv \frac{\hat{c}_p \mu}{k} \quad (\text{fluid properties})$$

Now, do the experiments.

35
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Complex Heat Transfer – Dimensional Analysis



Now, do the experiments.

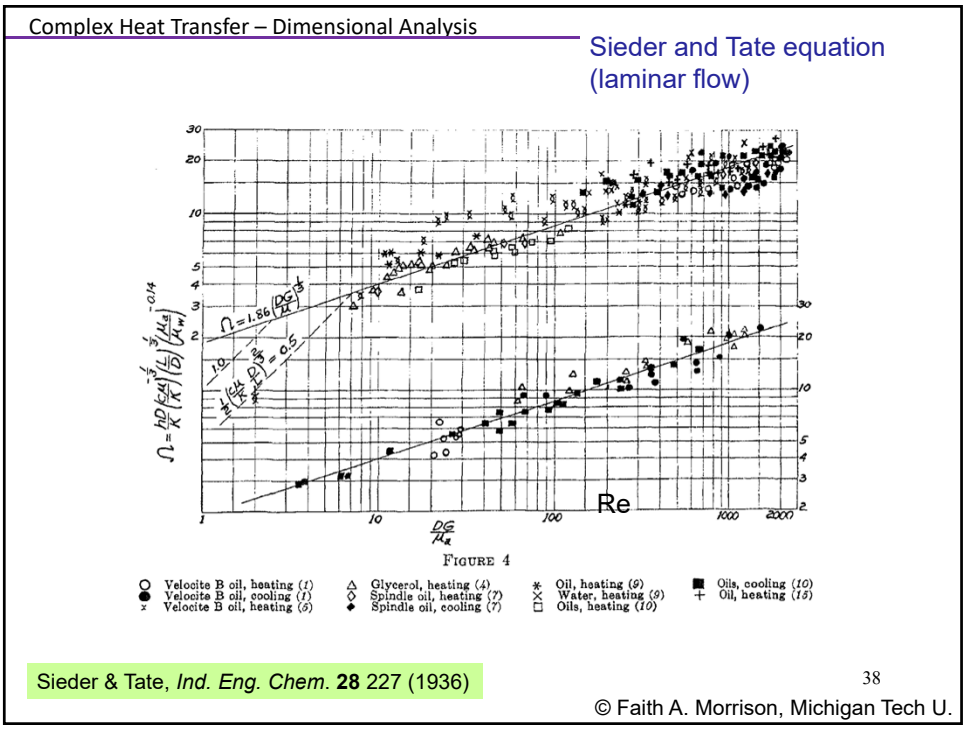
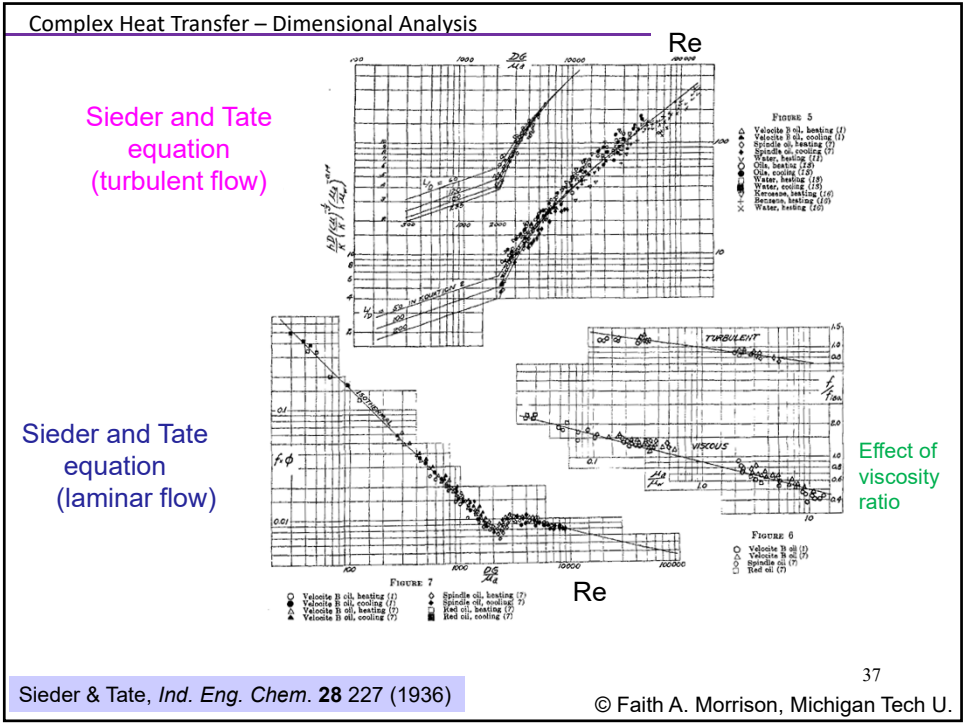
Forced Convection Heat Transfer

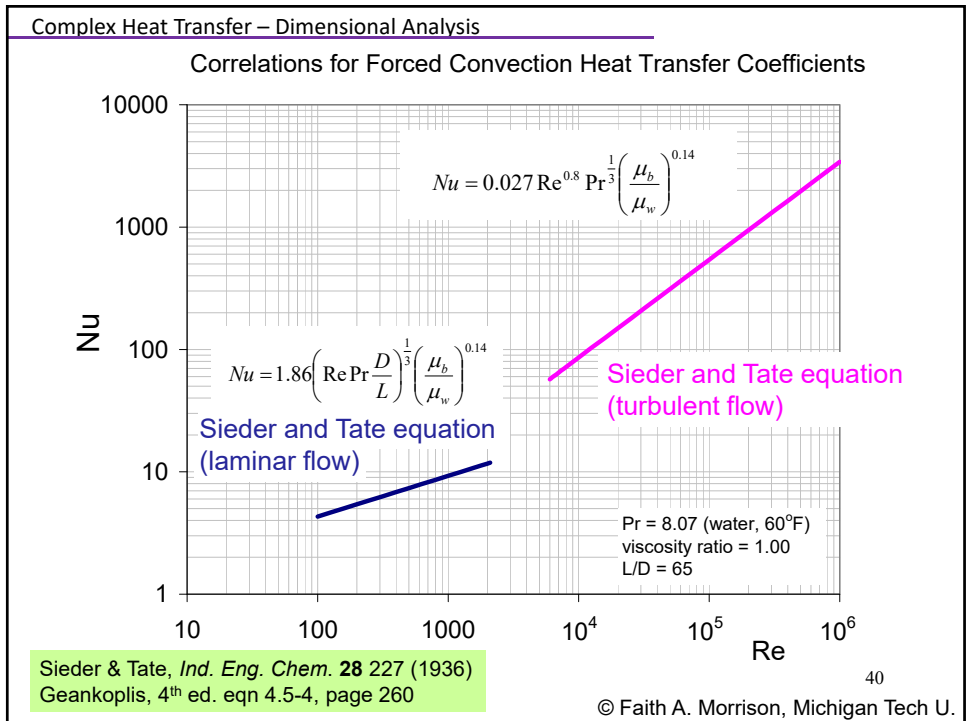
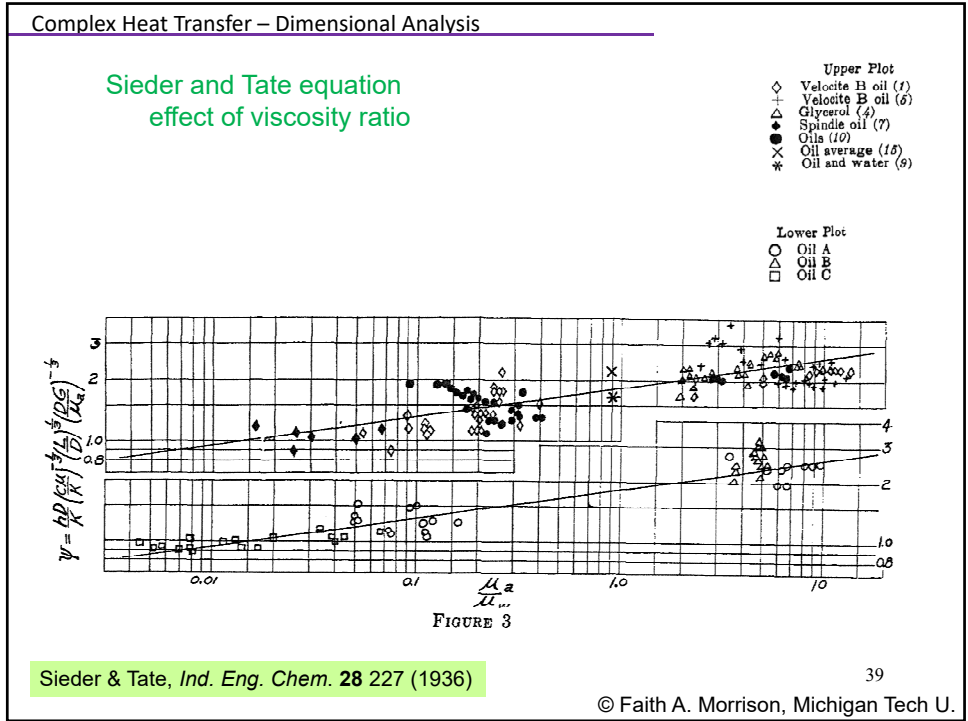
- Build apparatus (*several* actually, with different D, L)
- Run fluid through the inside (at different \underline{v} ; for different fluids ρ, μ, \hat{c}_p, k)
- Measure T_{bulk} on inside; T_{wall} on inside
- Measure rate of heat transfer, Q
- Calculate h : $|Q| = hA|T_{bulk} - T_{wall}|$
- Report h values in terms of dimensionless correlation:

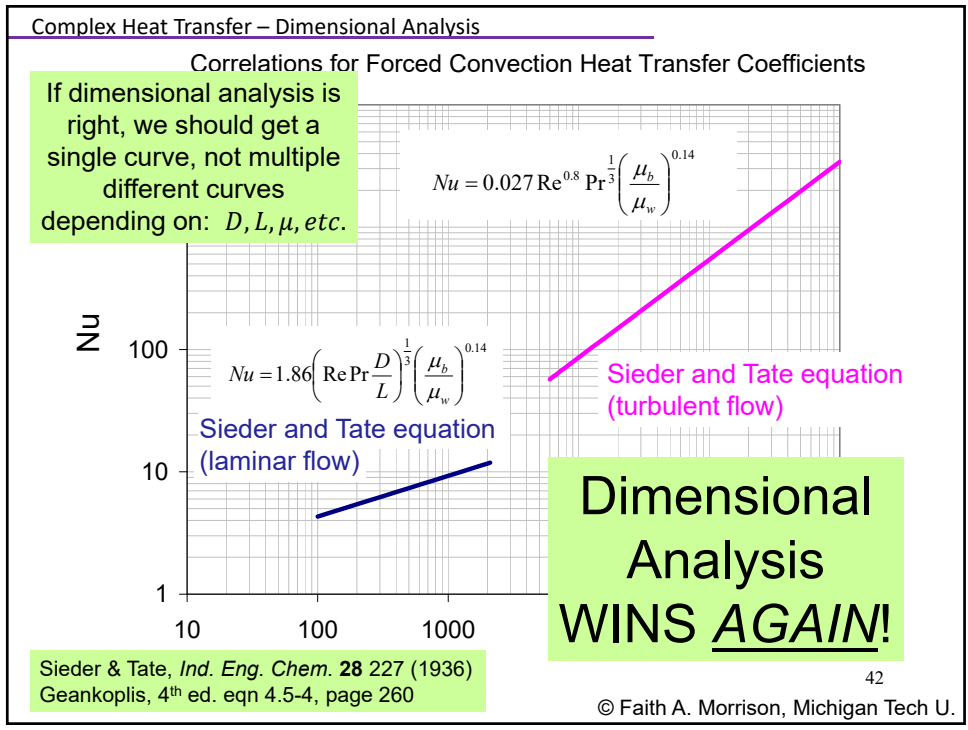
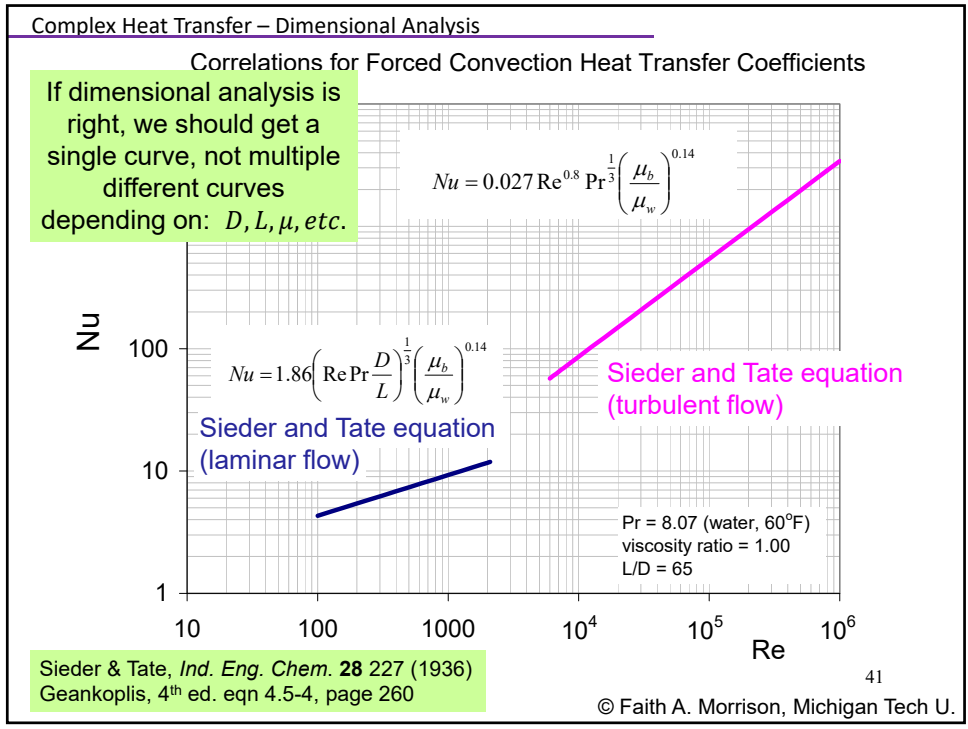
$$Nu = \frac{hD}{k} = f \left(Re, Pr, \frac{L}{D} \right)$$

It should only be a function of these dimensionless numbers (**if** our Dimensional Analysis is correct.....)

36
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Complex Heat Transfer – Dimensional Analysis

Heat Transfer in **Laminar** flow in pipes:
 data correlation for *forced convection heat transfer coefficients*

$$Nu_a = \frac{h_a D}{k} = 1.86 \left(Re Pr \frac{D}{L} \right)^{\frac{1}{3}} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Sieder-Tate equation (laminar flow)

Sieder & Tate, *Ind. Eng. Chem.* **28** 227 (1936)

the subscript “a” refers to *the type of average temperature* used in calculating the heat flow, q

$$q = h_a A \Delta T_a$$

$$\Delta T_a = \frac{(T_w - T_{bi}) + (T_w - T_{bo})}{2}$$

Geankoplis, 4th ed. eqn 4.5-4, page 260

$Re < 2100$, $(Re Pr \frac{D}{L}) > 100$, horizontal pipes; all physical properties evaluated at the mean temperature of the bulk fluid except μ_w which is evaluated at the (constant) wall temperature.

43

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Complex Heat Transfer – Dimensional Analysis

Heat Transfer in **Turbulent** flow in pipes:
 data correlation for *forced convection heat transfer coefficients*

$$Nu_{lm} = \frac{h_{lm} D}{k} = 0.027 Re^{0.8} Pr^{\frac{1}{3}} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Sieder-Tate equation (turbulent flow)

Sieder & Tate, *Ind. Eng. Chem.* **28** 227 (1936)

the subscript “lm” refers to *the type of average temperature* used in calculating the heat flow, q

$$q = h_{lm} A \Delta T_{lm}$$

$$\Delta T_{lm} = \frac{\Delta T_{w-bi} - \Delta T_{w-bo}}{\ln \left(\frac{\Delta T_{w-bi}}{\Delta T_{w-bo}} \right)}$$

Geankoplis, 4th ed. section 4.5

44

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Complex Heat Transfer – Dimensional Analysis

Forced convection
Heat Transfer in Laminar flow in pipes

$$Nu_a = \frac{h_a D}{k} = 1.86 \left(\text{RePr} \frac{D}{L} \right)^{\frac{1}{3}} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Sieder-Tate equation (laminar flow)

$$q = h_a A \Delta T_a$$

$$\Delta T_a = \frac{(T_w - T_{bi}) + (T_w - T_{bo})}{2}$$

Forced convection
Heat Transfer in Turbulent flow in pipes

$$Nu_{lm} = \frac{h_{lm} D}{k} = 0.027 \text{Re}^{0.8} \text{Pr}^{\frac{1}{3}} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Sieder-Tate equation (turbulent flow)

$$q = h_{lm} A \Delta T_{lm}$$

$$\Delta T_{lm} = \frac{\Delta T_{w-bi} - \Delta T_{w-bo}}{\ln \left(\frac{\Delta T_{w-bo}}{\Delta T_{w-bi}} \right)}$$

Physical Properties
(except μ_w) evaluated at:
 $\frac{T_{b,in} + T_{b,out}}{2}$

May have to be estimated

Bulk mean temperature

Fine print matters!

45

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Complex Heat Transfer – Dimensional Analysis

Forced convection Heat Transfer in Laminar flow in pipes

? In our dimensional analysis, we assumed constant ρ , k , μ , etc. Therefore we did not predict a viscosity-temperature dependence. If viscosity is not assumed constant, the dimensionless group shown below is predicted to appear in correlations.

$$Nu_a = \frac{h_a D}{k} = 1.86 \left(\text{RePr} \frac{D}{L} \right)^{\frac{1}{3}} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Sieder-Tate equation (laminar flow)

(reminiscent of pipe wall roughness; needed to modify dimensional analysis to correlate on roughness)

Sieder & Tate, Ind. Eng. Chem. 28 227 (1936)

46

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Complex Heat Transfer – Dimensional Analysis

Viscous fluids with large ΔT

heating

lower viscosity fluid layer speeds flow near the wall => higher h

$\mu_b > \mu_w$

cooling

higher viscosity fluid layer retards flow near the wall => lower h

$\mu_b < \mu_w$

empirical result:

$$\left(\frac{\mu_b}{\mu_w}\right)^{0.14}$$

ref: McCabe, Smith, Harriott, 5th ed, p339
Sieder & Tate, Ind. Eng. Chem. 28 227 (1936)

47
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Complex Heat Transfer – Dimensional Analysis

Forced convection
Heat Transfer in Laminar flow in pipes

$$Nu_a = \frac{h_a D}{k} = 1.86 \left(\text{RePr} \frac{D}{L} \right)^{\frac{1}{3}} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Why does $\frac{L}{D}$ appear in laminar flow correlations and not in the turbulent flow correlations?

LAMINAR

Less lateral mixing in laminar flow means more variation in $h(x)$.

48
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Complex Heat Transfer – Dimensional Analysis

Forced convection
Heat Transfer in Turbulent flow in pipes

$$Nu_{tm} = \frac{h_{tm} D}{k} = 0.027 Re^{0.8} Pr^{\frac{1}{3}} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Sieder & Tate eqn, turbulent flow

TURBULENT

In turbulent flow, good lateral mixing reduces the variation in h along the pipe length.

49
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(Exam 4 2016)


Example:

Water flows at 0.0522 kg/s (turbulent) in the inside of a double pipe heat exchanger (inside steel pipe, inner diameter = 0.545 inches , length unknown, physical properties given on page 1); the water enters at 30.0°C and exits at 65.6°C . In the shell of the heat exchanger, steam condenses at an unknown saturation pressure. What is the heat transfer coefficient, h_{lm} (based on log mean temperature driving force) in the water flowing in the pipe? You may neglect the effect on heat-transfer coefficient of the temperature-dependence of viscosity. Please give your answer in $\text{W/m}^2\text{K}$.

Physical properties of steel:
 thermal conductivity = 16.3 W/mK
 heat capacity = 0.49 kJ/kg K
 density = 8050 kg/m^3

50
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This is shown to work for *internal flow* (Seider-Tate).

Complex Heat Transfer – Dimensional Analysis 

Now, do the experiments.

Forced Convection Heat Transfer

- Build apparatus (*several* actually, with different D, L)
- Run fluid through the inside (at different v ; for different fluids ρ, μ, \hat{c}_p, k)
- Measure T_{bulk} on inside; T_{wall} on inside
- Measure rate of heat transfer, Q
- Calculate h : $|Q| = hA|T_{bulk} - T_{wall}|$
- Report h values in terms of dimensionless correlation:

$$Nu = \frac{hD}{k} = f\left(Re, Pr, \frac{L}{D}\right)$$

It should only be a function of these dimensionless numbers (*if* our Dimensional Analysis is correct.....)

It also works for **external** flows (flows past spheres, plates, cylinders, banks of cylinders (see Geankoplis, section 4.6))

Complex Heat Transfer – Dimensional Analysis

Heat Transfer in **External** flow around objects:
data correlation for forced convection heat transfer coefficients

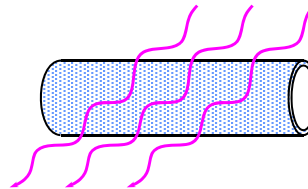
$$Nu = \frac{hD_{ext}}{k} = CRe^m Pr^{\frac{1}{3}}$$

$q = hA(T_w - T_b)$ D_{ext} = external diameter

External flow, perpendicular to cylinder

Geankoplis, 4th ed. section 4.6, p272

Re	m	C
1 – 4	0.330	0.989
4 – 40	0.385	0.911
40 – 4,000	0.466	0.683
4,000 – 4×10^4	0.618	0.193
4×10^4 – 2.5×10^5	0.805	0.0266



Dimensional Analysis

These numbers tell us about the relative importance of the terms they precede.

Dimensionless numbers from the Equations of Change (microscopic balances)

momentum

Non-dimensional Navier-Stokes Equation

$$\left(\frac{\partial v_z^*}{\partial t^*} + \underline{v}^* \cdot \nabla^* v_z^* \right) = - \frac{\partial P^*}{\partial z^*} + \frac{1}{\text{Re}} (\nabla^{*2} v_z^*) + \frac{1}{\text{Fr}} g^*$$

energy

Non-dimensional Energy Equation

$$\left(\frac{\partial T^*}{\partial t^*} + \underline{v}^* \cdot \nabla^* T^* \right) = \frac{1}{\text{RePr}} (\nabla^{*2} T^*) + S^*$$

mass

Non-dimensional Continuity Equation (species A)

$$\left(\frac{\partial x_A^*}{\partial t^*} + \underline{v}^* \cdot \nabla^* x_A^* \right) = \frac{1}{\text{ReSc}} (\nabla^{*2} x_A^*)$$

Re – Reynolds
Fr – Froude

Pe – Péclet_h = RePr
Pr – Prandtl

Pe – Péclet_m = ReSc
Sc – Schmidt

ref: BSL1, p581, 644 53

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Dimensionless Numbers

Dimensionless numbers from the Equations of Change

Re – Reynolds = $\frac{\rho V D}{\mu} = \frac{V D}{\nu}$

Fr – Froude = $\frac{V^2}{g D}$

Pe – Péclet_h = **RePr** = $\frac{\hat{c}_p \rho V D}{k} = \frac{V D}{\alpha}$

Pe – Péclet_m = **ReSc** = $\frac{V D}{D_{AB}}$

Pr – Prandtl = $\frac{\hat{c}_p \mu}{k} = \frac{\nu}{\alpha}$

Sc – Schmidt = **LePr** = $\frac{\mu}{\rho D_{AB}} = \frac{\nu}{D_{AB}}$

Le – Lewis = $\frac{\alpha}{D_{AB}}$

These numbers tell us about the **relative importance of the terms** they precede in the microscopic balances (*scenario properties*).

These numbers compare the **magnitudes of the diffusive transport coefficients** ν, α, D_{AB} (*material properties*).

54

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Dimensionless Numbers

Dimensionless numbers from the **Equations of Change**

Re – Reynolds = $\frac{\rho V D}{\mu} = \frac{V D}{\nu}$

Fr – Froude = $\frac{V^2}{g D}$

Pe – Péclet_h = **RePr** = $\frac{\hat{C}_p \rho V D}{k} = \frac{V D}{\alpha}$

Pe – Péclet_m = **ReSc** = $\frac{V D}{D_{AB}}$

Pr – Prandtl = $\frac{\hat{C}_p \mu}{k} = \frac{\nu}{\alpha}$

Sc – Schmidt = **LePr** = $\frac{\mu}{\rho D_{AB}} = \frac{\nu}{D_{AB}}$

Le – Lewis = $\frac{\alpha}{D_{AB}}$

These numbers tell us about the **relative importance of the terms** they precede in the microscopic balances (*scenario properties*).

These numbers compare the **magnitudes of the diffusive transport coefficients** ν, α, D_{AB} (*material properties*).

Transport coefficients

55

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Dimensional Analysis

Dimensionless numbers from the **Engineering Quantities of Interest**

momentum

Dimensionless Force on the Wall (Drag)

$$f = \frac{1}{\pi L \text{Re}} \int_0^{\frac{L}{2}} \int_0^{2\pi} \left(\frac{\partial v_z^*}{\partial r^*} \right) \Big|_{r^*=\frac{1}{2}} d\theta dz^*$$

energy

Newton's Law of Cooling

$$\text{Nu} = \frac{1}{2\pi L / D} \int_0^{\frac{L}{2}} \int_0^{2\pi} \left(\frac{\partial T^*}{\partial r^*} \right) \Big|_{r^*=1/2} dz^* d\theta$$

mass xfer

Dimensionless Mass Transfer Coefficient

$$\text{Sh} = \frac{1}{2\pi L} \int_0^{\frac{L}{2}} \int_0^{2\pi} \left(-\frac{\partial x_A}{\partial r^*} \right) \Big|_{r^*=1/2} d\theta dz^*$$

These numbers are defined to help us build **transport data correlations** based on the fewest number of grouped (dimensionless) variables (*scenario property*).

f – Friction Factor (Fanning)

L/D – Aspect Ratio

$$f = \frac{\mathcal{F}_{drag}}{\left(\frac{1}{2}\rho V^2\right) A_c}$$

Nu – Nusselt

L/D – Aspect Ratio

$$\text{Nu} = \frac{hD}{k}$$

Sh – Sherwood

L/D – Aspect Ratio

$$\text{Sh} = \frac{k_m D}{D_{AB}}$$

56

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momentum
energy
mass

Dimensionless Numbers

$\text{Re} - \text{Reynolds} = \frac{\rho V D}{\mu} = \frac{V D}{\nu}$	<div style="background-color: #f08080; padding: 5px; border: 1px solid black;"> These numbers from the governing equations tell us about the relative importance of the terms they precede in the </div>
$\text{Fr} - \text{Froude} = \frac{v^2}{g D}$	
$\text{Pe} - \text{Péclet}_h = \text{RePr} = \frac{\hat{c}_p \rho V D}{\kappa} = \frac{V D}{\alpha}$	
$\text{Pe} - \text{Péclet}_m = \text{ReSc} = \frac{V D}{D_{AB}}$	<div style="background-color: #00b050; color: white; padding: 5px; border: 1px solid black;"> These numbers compare the magnitudes of the diffusive transport coefficients ν, α, D_{AB} </div>
$\text{Pr} - \text{Prandtl} = \frac{\hat{c}_p \mu}{\kappa} = \frac{\nu}{\alpha}$	<div style="background-color: #0000ff; color: white; padding: 5px; border: 1px solid black;"> These numbers are defined to help us build transport data correlations based on the fewest number of grouped (dimensionless) variables (scenario) </div>
$\text{Sc} - \text{Schmidt} = \text{LePr} = \frac{\mu}{\rho D_{AB}} = \frac{\nu}{D_{AB}}$	
$\text{Le} - \text{Lewis} = \frac{\alpha}{D_{AB}}$	
$f - \text{Friction Factor} =$	<div style="font-size: small; text-align: right;">© Faith A. Morrison, Michigan Tech U.</div>

Complex Heat Transfer – Dimensional Analysis

Complex Heat transfer Problems to Solve:

- ✓ • Forced convection heat transfer from fluid to wall
Solution: ?
- ➔ • Natural convection heat transfer from fluid to wall
Solution: ?
- Radiation heat transfer from solid to fluid
Solution: ?

We started with a forced-convection pipe problem, did dimensional analysis, and found the dimensionless numbers.

To do a situation with different physics, we must start with a different starting problem.

58
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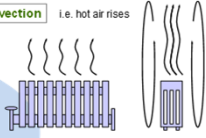
Next:

CM3110
Transport/Unit Ops I
Part II: Heat Transfer

Michigan Tech

Free Convection i.e. hot air rises

Complex Heat Transfer – Dimensional Analysis
(Natural convection heat transfer)



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