

Types of Heat Transfer

$$\frac{q_x}{A} = -k \frac{dT}{dx} \quad \bullet \text{conduction (Fourier's Law)}$$

$$\underline{v} \cdot \nabla T \quad \bullet \text{forced convection (due to flow)}$$

$$S \quad \bullet \text{source terms}$$

$$\frac{Dv_z^*}{Dt} = (\nabla^2 v_z^*) + GrT^* \quad \bullet \text{free convection (fluid motion due to density variations brought on by temperature differences)}$$

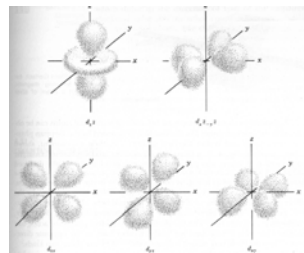
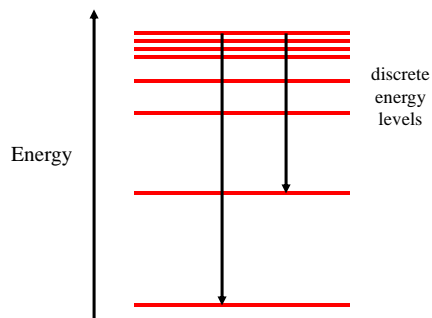
$$\Delta H_{vap} \quad \bullet \text{heat transfer with phase change (e.g. condensing fluids)}$$

last subject in the course { **•radiation**

© Faith A. Morrison, Michigan Tech U.

Heat transfer due to radiation

- in atoms and molecules electrons can exist in multiple, discrete energy states
- transfers between energy states are accompanied by an emission of radiation



Sienko and Plane, Chemistry: Principles and Applications, McGraw Hill, 1979

Quantum Mechanics

© Faith A. Morrison, Michigan Tech U.

Radiation versus Conduction and Convection

Continuum view

- Conduction is caused by macroscopic temperature gradients
- Convection is caused by macroscopic flow
- Radiation? **NO CONTINUUM EXPLANATION**

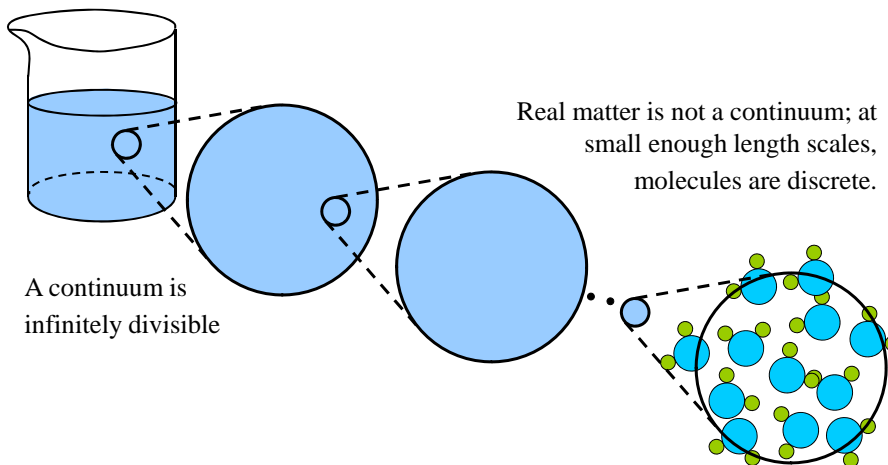
Molecular view

- Conduction? }
- Convection? }
- Radiation is caused by changes in electron energy states in molecules and atoms

There is, of course, a molecular explanation of these effects, since we know that matter is made of atoms and molecules

© Faith A. Morrison, Michigan Tech U.

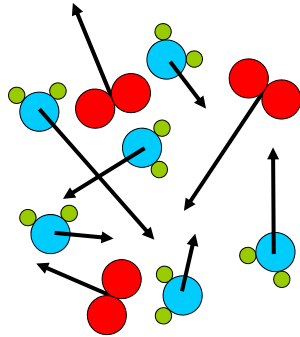
Continuum versus Molecular description of matter



© Faith A. Morrison, Michigan Tech U.

Individual molecules carry:

- chemical identity
- macroscopic velocity (speed and direction)
- internal energy (Brownian velocity)



When they undergo Brownian motion within an inhomogeneous mixture, they cause:

- diffusion (mass transport)
- exchange of momentum (momentum transport)
- conduction (energy transport)

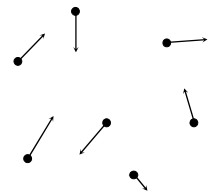
© Faith A. Morrison, Michigan Tech U.

Kinetic Theory J. C. Maxwell, L. Boltzmann, 1860

- Molecules are in constant motion (Brownian motion)
- Temperature is related to $E_{k,av}$ of the molecules

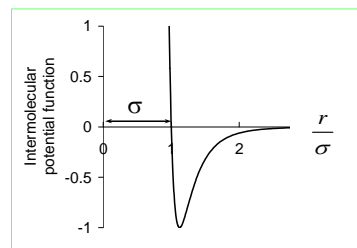
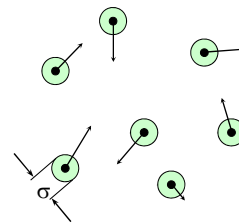
Simplest model

- no particle volume
- no intermolecular forces



More realistic model

- finite particle volume
- intermolecular forces



© Faith A. Morrison, Michigan Tech U.

Kinetic Theory

Is based on Brownian motion (molecules in constant motion proportional to their temperature)

Predicts that properties that are carried by individual molecules (chemical identity, momentum, average kinetic energy) will be transported **DOWN** gradients in these quantities.

==> Transport laws are due to Brownian motion

Heat Transfer by Radiation

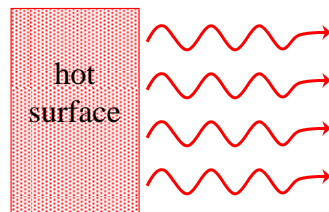
Is due to the release of energy stored in molecules that is **NOT** related to average kinetic energy (temperature), but rather to the population of excited states.

==> Radiation is **NOT** a Brownian effect

© Faith A. Morrison, Michigan Tech U.

Radiation

- does not require a medium to transfer energy (works in a vacuum)
- travels at the speed of light, $c = 3 \times 10^{10}$ cm/s
- travels as a *wave*; differs from x-rays, light, only by wavelength, λ
- radiation is important when temperatures are high



examples:

- the sun
- home radiator
- hot walls in vacuum oven
- heat exchanger walls when ΔT is high and a vapor film has formed

$$\frac{q}{A} \propto T^4$$

© Faith A. Morrison, Michigan Tech U.

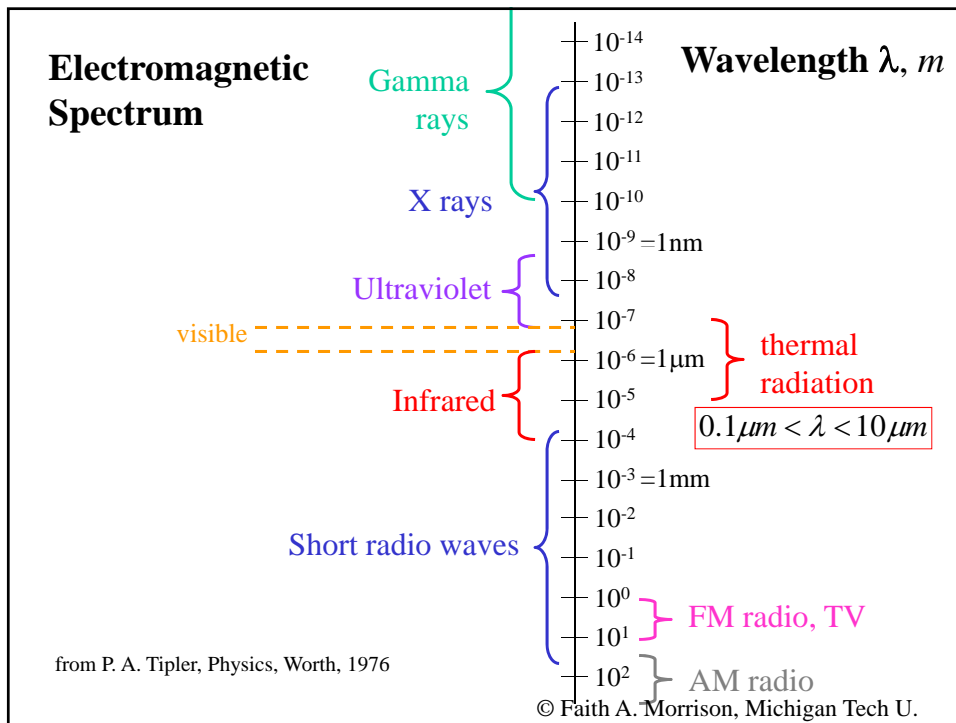
Why does radiation flux scale with temperature, which is related to average kinetic energy?

As a molecule gains energy, it both speeds up (increases average kinetic energy) and increases its population of excited states.

The increase in **average kinetic energy** is reflected in temperature (directly proportional).

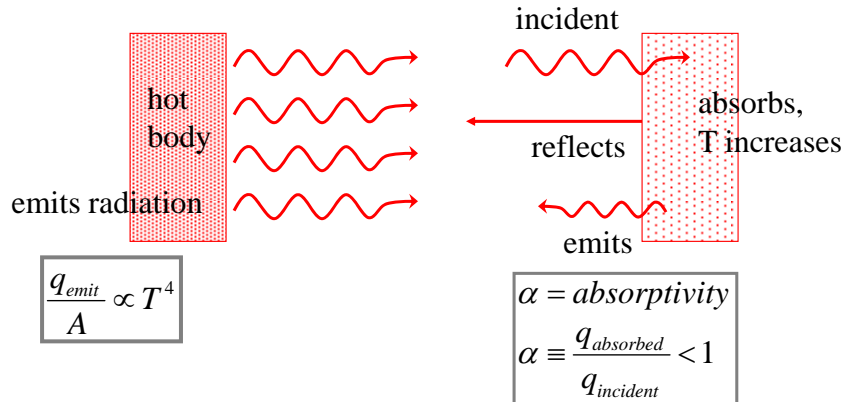
The increase in number of electrons in **excited states** is reflected in increased radiation flux. Electrons enter excited states in proportion to T^4 .

© Faith A. Morrison, Michigan Tech U.



What causes energy transfer by radiation?

- energy hits surface
- pushes some molecules into an excited state
- when the molecules/atoms relax from the excited state, they emit radiation



© Faith A. Morrison, Michigan Tech U.

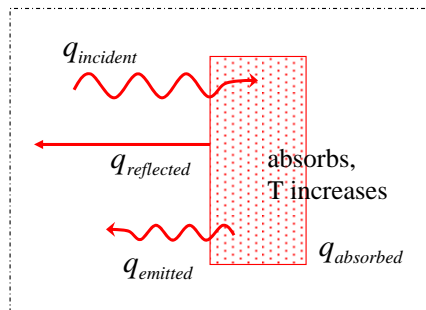
Absorption

$$\alpha = \text{absorptivity}$$

$$\alpha \equiv \frac{q_{absorbed}}{q_{incident}} < 1$$

In general, α is a function of wavelength

$$\alpha = \alpha(\lambda)$$



gray body: a body for which α is constant (does not depend on λ)

black body: a body for which $\alpha = 1$, i.e. absorbs all incident radiation

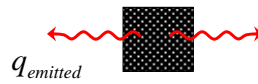
© Faith A. Morrison, Michigan Tech U.

Emission

$\epsilon = \text{emissivity}$

$$\epsilon \equiv \frac{q_{\text{emitted}}}{q_{\text{emitted,black body}}} < 1$$

gray body: a body for which α is constant
black body: a body for which $\alpha = 1$



$$\alpha = \text{absorptivity}$$

$$\alpha \equiv \frac{q_{\text{absorbed}}}{q_{\text{incident}}} < 1$$

Kirchhoff's Law: emissivity equals absorptivity at the same temperature

$$\alpha = \epsilon$$

true for
black and
non-black
solid
surfaces

the fraction of energy absorbed by a material = the relative amount of energy emitted from that material compared to a black body

© Faith A. Morrison, Michigan Tech U.

$\epsilon = \text{emissivity}$

$$\epsilon \equiv \frac{q_{\text{emitted}}}{q_{\text{emitted,black body}}} < 1$$

Black Bodies

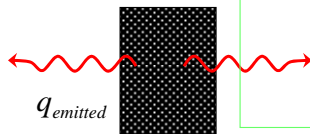
Stefan-Boltzmann Law: the amount of energy emitted by a black body is proportional to T^4

$$\frac{q_{\text{emitted,black body}}}{A} = \sigma T^4$$

absolute temperature

$$\sigma = 0.1712 \times 10^{-8} \frac{\text{BTU}}{\text{h ft}^2 \text{R}^4}$$

$$= 5.676 \times 10^{-8} \frac{\text{W}}{\text{m}^2 \text{K}^4}$$



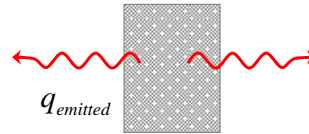
© Faith A. Morrison, Michigan Tech U.

Non-Black Bodies

$\varepsilon = \text{emissivity}$

$$\varepsilon \equiv \frac{q_{\text{emitted}}}{q_{\text{emitted, black body}}}$$

$$\frac{q_{\text{emitted, non-black body}}}{A} = \varepsilon q_{\text{emitted, black body}} \\ = \varepsilon \sigma T^4$$



Stefan-Boltzmann:

$$\frac{q_{\text{emitted, black body}}}{A} = \sigma T^4$$

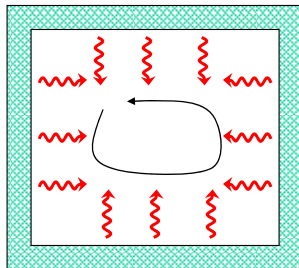
Energy emitted by a non-black body

$$\frac{q_{\text{emitted, non-black body}}}{A} = \varepsilon \sigma T^4$$

© Faith A. Morrison, Michigan Tech U.

How does this relate to chemical engineering?

Consider a furnace with an internal blower:



There is heat transfer due to convection:

$$q_{\text{convection}} = h_{\text{conv}} A (T_s - T_b)$$

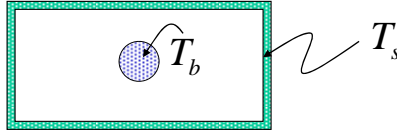
There is also heat transfer due to radiation:

$$q_{\text{radiation}} = h_{\text{rad}} A (T_s - T_b)$$

$$q_{\text{total}} = q_{\text{conv}} + q_{\text{rad}}$$

© Faith A. Morrison, Michigan Tech U.

Where do we get h_{rad} ?



object in furnace: $q_{emitted, non-blackbody} = A\varepsilon|_{T_b} \sigma T_b^4$ using Kirchhoff's law

$q_{absorbed} = \alpha|_{T_s} A \sigma T_s^4 = A\varepsilon|_{T_s} \sigma T_s^4$

energy emitted by walls, which are acting as a black body

net energy absorbed:

$$q_{transferred\ to\ body} = A\varepsilon|_{T_s} \sigma (T_s^4 - T_b^4)$$

emissivity at T_s

assuming $\varepsilon|_{T_s} \approx \varepsilon|_{T_b}$

© Faith A. Morrison, Michigan Tech U.

Finally, calculate h_{rad}

net energy absorbed:

$$q_{transferred\ to\ body} = A\varepsilon|_{T_s} \sigma (T_s^4 - T_b^4)$$

assuming $\varepsilon|_{T_s} \approx \varepsilon|_{T_b}$

equating with expression for h: $A\varepsilon|_{T_s} \sigma (T_s^4 - T_b^4) = h_{rad} A(T_s - T_b)$

$$h_{rad} = \frac{\sigma \varepsilon|_{T_s} (T_s^4 - T_b^4)}{T_s - T_b}$$

Geankoplis 4th ed., eqn 4.10-10 p304

© Faith A. Morrison, Michigan Tech U.

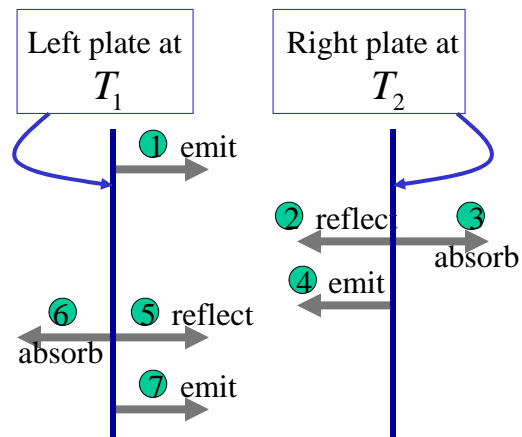
Example: Geankoplis 4.10-3

A horizontal oxidized steel pipe carrying steam and having an OD of 0.1683m has a surface temperature of 374.9 K and is exposed to air at 297.1 K in a large enclosure. Calculate the heat loss for 0.305 m of pipe from natural convection plus radiation. For the steel pipe, use an emissivity of 0.79.

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer Between Two Infinite Plates

Consider a quantity of radiation energy that is emitted from surface 1.



See: Geankoplis, section 4.11B

Also: Bird, Stewart, and Lightfoot, "Transport Phenomena" 1960 Wiley PP446-448

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer
Between Two Infinite
Plates

**First round –
surface 2**

Quantity of energy **incident** at surface 2: $\frac{q_{1-2}}{A} = \epsilon_1 \sigma T_1^4$

Quantity of energy **absorbed** at surface 2: $\alpha_2 \left(\frac{q_{1-2}}{A} \right) A = \epsilon_2 (\epsilon_1 \sigma T_1^4) A$
 $\alpha_2 = \epsilon_2$

Quantity of energy **reflected** from surface 2: $(1 - \epsilon_2) (\epsilon_1 A \sigma T_1^4)$
 fraction reflected incident energy

↑
This energy goes back to surface 1.

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer
Between Two Infinite
Plates

**Second round –
surface 1**

Quantity of energy **absorbed** at surface 1 (second round): $\epsilon_1 \left[(1 - \epsilon_2) (\epsilon_1 A \sigma T_1^4) \right]$
 fraction absorbed incident energy

Quantity of energy **reflected** from surface 1 (second round): $(1 - \epsilon_1) \left[(1 - \epsilon_2) (\epsilon_1 A \sigma T_1^4) \right]$
 fraction reflected incident energy

© Faith A. Morrison, Michigan Tech U.

Third round – surface

Radiation Heat Transfer
Between Two Infinite
Plates

Quantity of energy
absorbed at surface 2
(third round):

$$\underbrace{\varepsilon_2}_{\text{fraction absorbed}} \underbrace{\left[(1 - \varepsilon_1)(1 - \varepsilon_2) (\varepsilon_1 A \sigma T_1^4) \right]}_{\text{incident energy}}$$

Quantity of energy
reflected from surface 2
(third round):

$$\underbrace{(1 - \varepsilon_2)}_{\text{fraction reflected}} \underbrace{\left[(1 - \varepsilon_1)(1 - \varepsilon_2) (\varepsilon_1 A \sigma T_1^4) \right]}_{\text{incident energy}}$$

There is a pattern.

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer
Between Two Infinite
Plates

Now, calculate the radiation energy
going from surface 1 to surface 2:

Later, calculate energy from
2 to 1; then subtract to obtain
net energy transferred.

$$q_{1-2} = \left(\begin{array}{c} \text{energy from} \\ 1 \rightarrow 2 \end{array} \right) = \sum \left(\begin{array}{c} \text{energy absorbed} \\ \text{at surface 2} \end{array} \right)$$

$$= \varepsilon_2 (\varepsilon_1 A \sigma T_1^4)$$

$$+ \varepsilon_2 (1 - \varepsilon_1)(1 - \varepsilon_2) (\varepsilon_1 A \sigma T_1^4)$$

$$+ \varepsilon_2 (1 - \varepsilon_1)^2 (1 - \varepsilon_2)^2 (\varepsilon_1 A \sigma T_1^4)$$

$$\dots + \varepsilon_2 (1 - \varepsilon_1)^n (1 - \varepsilon_2)^n (\varepsilon_1 A \sigma T_1^4) + \dots$$

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer
Between Two Infinite
Plates

Radiation energy going from
surface 1 to surface 2:

$$q_{1-2} = \varepsilon_1 \varepsilon_2 A \sigma T_1^4 \sum_{n=0}^{\infty} (1 - \varepsilon_1)^n (1 - \varepsilon_2)^n$$

How can we calculate $\sum_{n=0}^{\infty} x^n$?

Answer: $S = \frac{1}{1-x}$

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer
Between Two Infinite
Plates

Radiation energy going from
surface 1 to surface 2:

$$q_{1-2} = \frac{\varepsilon_1 \varepsilon_2 A \sigma T_1^4}{1 - [(1 - \varepsilon_1)(1 - \varepsilon_2)]}$$

$$= \frac{\varepsilon_1 \varepsilon_2 A \sigma T_1^4}{1 - [1 - \varepsilon_1 - \varepsilon_2 + \varepsilon_1 \varepsilon_2]} = \frac{\varepsilon_1 \varepsilon_2 A \sigma T_1^4}{\varepsilon_1 + \varepsilon_2 - \varepsilon_1 \varepsilon_2}$$

$$\frac{q_{1-2}}{A} = \frac{\sigma T_1^4}{\frac{1}{\varepsilon_1} + \frac{1}{\varepsilon_2} - 1}$$

Final Result

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer Between Two Infinite Plates

Radiation energy going from
surface **1** to surface **2**:

$$\frac{q_{1-2}}{A} = \frac{\sigma T_1^4}{\frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} - 1}$$

Radiation energy going from
surface **2** to surface **1**:

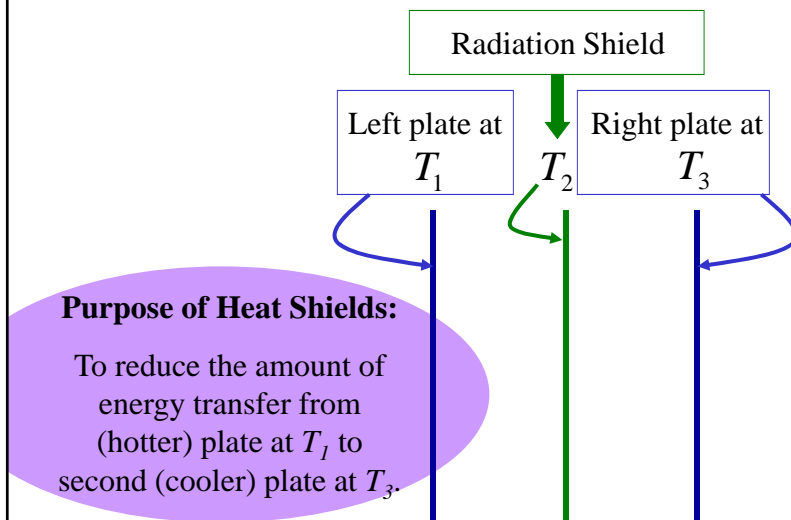
$$\frac{q_{2-1}}{A} = \frac{\sigma T_2^4}{\frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} - 1}$$

NET Radiation energy going
from surface 1 to surface 2:

$$\frac{q_{1-2} - q_{2-1}}{A} = \frac{\sigma (T_1^4 - T_2^4)}{\left(\frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} - 1 \right)}$$

© Faith A. Morrison, Michigan Tech U.

Radiation Shields



Purpose of Heat Shields:

To reduce the amount of energy transfer from (hotter) plate at T_1 to second (cooler) plate at T_3 .

Note:

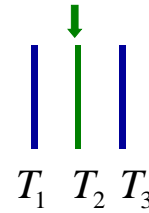
$$q_{net,1 \rightarrow 2} = q_{net,2 \rightarrow 3} = q$$

© Faith A. Morrison, Michigan Tech U.

Analysis of Radiation Shields

We will assume that the emissivity is the same for all surfaces.

Radiation Shield



$$\frac{q_{net,1 \rightarrow 2}}{A} = \frac{\sigma(T_1^4 - T_2^4)}{\left(\frac{1}{\varepsilon} + \frac{1}{\varepsilon} - 1\right)}$$

$$\frac{q_{net,2 \rightarrow 3}}{A} = \frac{\sigma(T_2^4 - T_3^4)}{\left(\frac{1}{\varepsilon} + \frac{1}{\varepsilon} - 1\right)}$$

Now we eliminate T_2 between these equations.

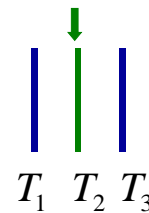
Note:

$$q_{net,1 \rightarrow 2} = q_{net,2 \rightarrow 3} = q$$

© Faith A. Morrison, Michigan Tech U.

Analysis of Radiation Shields

Radiation Shield



$$\frac{q}{A} = \frac{\sigma(T_1^4 - T_2^4)}{\left(\frac{2}{\varepsilon} - 1\right)} \quad \frac{q}{A} = \frac{\sigma(T_2^4 - T_3^4)}{\left(\frac{2}{\varepsilon} - 1\right)}$$

$$T_2^4 = \frac{q}{\sigma A} \left(\frac{2}{\varepsilon} - 1\right) + T_3^4$$

$$\frac{q}{\sigma A} \left(\frac{2}{\varepsilon} - 1\right) = T_1^4 - \frac{q}{\sigma A} \left(\frac{2}{\varepsilon} - 1\right) - T_3^4$$

$$\frac{2q}{\sigma A} \left(\frac{2}{\varepsilon} - 1\right) = T_1^4 - T_3^4$$

$$\frac{q}{A} = \left(\frac{1}{2}\right) \frac{\sigma(T_1^4 - T_3^4)}{(2/\varepsilon - 1)}$$

© Faith A. Morrison, Michigan Tech U.

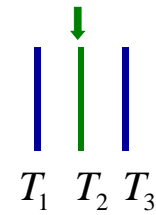
Analysis of Radiation Shields

**1 Heat
Shield**

$$\frac{q}{A} = \left(\frac{1}{2} \right) \frac{\sigma(T_1^4 - T_3^4)}{(2/\varepsilon - 1)}$$

With one heat shield present, q falls by half compared to no heat shield.

Radiation Shield



by the same analysis,

**N Heat
Shields**

$$\frac{q}{A} = \left(\frac{1}{N+1} \right) \frac{\sigma(T_1^4 - T_3^4)}{(2/\varepsilon - 1)}$$

With N heat shields present, q falls by a factor of $1/N$ compared to no heat shield.

© Faith A. Morrison, Michigan Tech U.