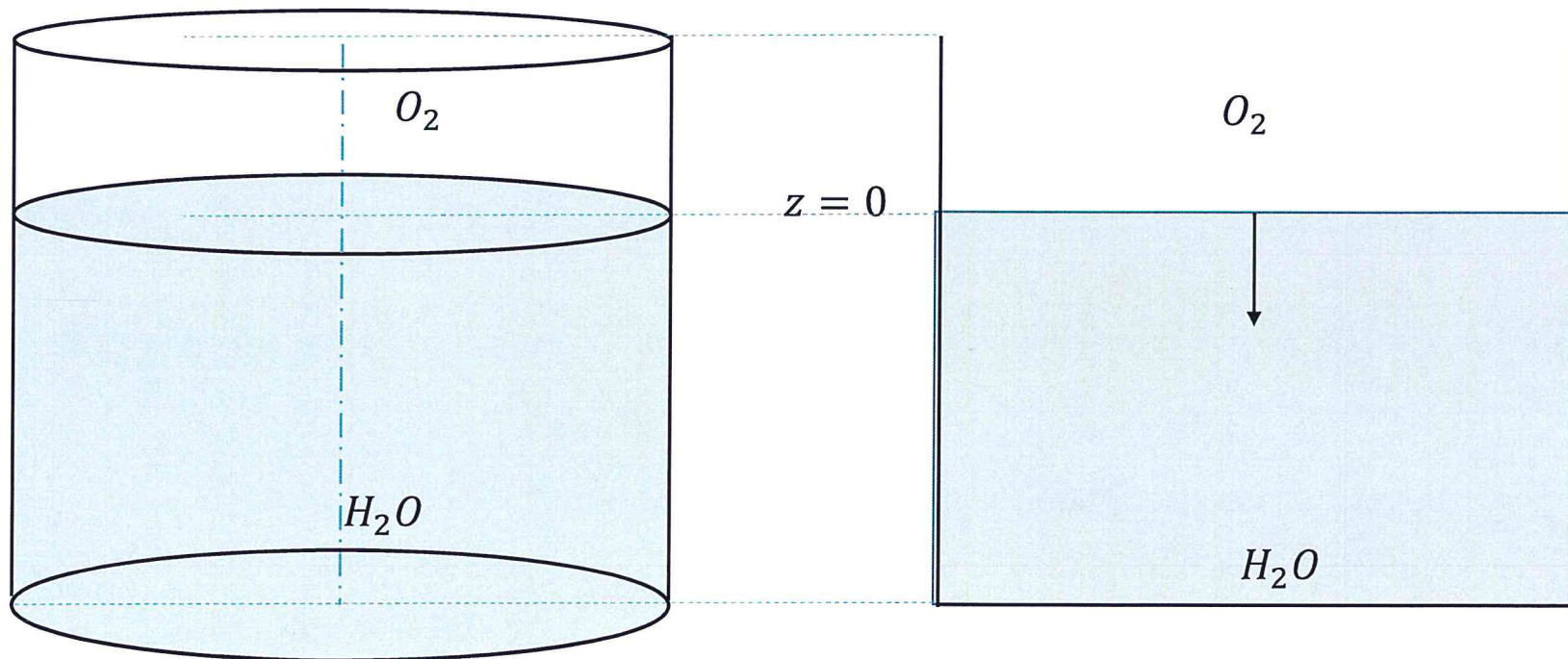


UNSTEADY STATE EXAMPLE



Unsteady State Mass Transport

Example: A very long, very large tank of water is suddenly exposed to oxygen atmosphere. Oxygen diffuses into the water. What is the concentration profile of the oxygen in the water as a function of time?



CARTESIAN

(2)

The **Equation of Species Mass Balance, constant ρD_{AB}** . For binary systems, and Fick's law has been incorporated. Good for dilute liquid solutions at constant temperature and pressure.

Microscopic species mass balance, constant thermal conductivity; Gibbs notation

$$\rho \left(\frac{\partial \omega_A}{\partial t} + \underline{v} \cdot \nabla \omega_A \right) = \rho D_{AB} \nabla^2 \omega_A + r_A$$

Microscopic species mass balance, constant thermal conductivity; Cartesian coordinates

$$\rho \left(\frac{\partial \omega_A}{\partial t} + v_x \frac{\partial \omega_A}{\partial x} + v_y \frac{\partial \omega_A}{\partial y} + v_z \frac{\partial \omega_A}{\partial z} \right) = \rho D_{AB} \left(\frac{\partial^2 \omega_A}{\partial x^2} + \frac{\partial^2 \omega_A}{\partial y^2} + \frac{\partial^2 \omega_A}{\partial z^2} \right) + r_A$$

Microscopic species mass balance, constant thermal conductivity; cylindrical coordinates

$$\rho \left(\frac{\partial \omega_A}{\partial t} + v_r \frac{\partial \omega_A}{\partial r} + \frac{v_\theta}{r} \frac{\partial \omega_A}{\partial \theta} + v_z \frac{\partial \omega_A}{\partial z} \right) = \rho D_{AB} \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \omega_A}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \omega_A}{\partial \theta^2} + \frac{\partial^2 \omega_A}{\partial z^2} \right) + r_A$$

Microscopic species mass balance, constant thermal conductivity; spherical coordinates

$$\rho \left(\frac{\partial \omega_A}{\partial t} + v_r \frac{\partial \omega_A}{\partial r} + \frac{v_\theta}{r} \frac{\partial \omega_A}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial \omega_A}{\partial \phi} \right) = \rho D_{AB} \left(\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \omega_A}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \omega_A}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \omega_A}{\partial \phi^2} \right) + r_A$$

In terms of Diffusivity, D_{AB} **no ω_A terms**

$$c x_A = c_A = \frac{1}{M_A} (\rho_A) = \frac{1}{M_A} (\rho \omega_A) \quad \left(\text{units: } c [=] \frac{\text{mol mix}}{\text{vol soln}} ; \rho [=] \frac{\text{mass mix}}{\text{vol soln}} ; c_A [=] \frac{\text{mol } A}{\text{vol soln}} ; \rho_A [=] \frac{\text{mass } A}{\text{vol soln}} \right)$$

$$\underline{J}_A \equiv \text{mass flux of species } A \text{ relative to a mixture's mass average velocity, } \underline{v} \quad \left(\text{units: } \underline{J}_A [=] \frac{\text{mass } A}{\text{area} \cdot \text{time}} \right)$$

$$= \rho_A (\underline{v}_A - \underline{v})$$

$\underline{J}_A + \underline{J}_B = 0$, i.e. these fluxes are measured relative to the mixture's center of mass

$\underline{n}_A \equiv \rho_A \underline{v}_A = \underline{J}_A + \rho_A \underline{v} =$ combined mass flux relative to stationary coordinates

$$\underline{n}_A + \underline{n}_B = \rho \underline{v}$$

$\underline{v}_A \equiv$ velocity of species A in a mixture, i.e. average velocity of all molecules of species A within a small volume

$\underline{v} = \omega_A \underline{v}_A + \omega_B \underline{v}_B \equiv$ mass average velocity; same velocity as in the microscopic momentum and energy balances

Reference: R. B. Bird, W. E. Stewart, and E. N. Lightfoot, *Transport Phenomena*, 2nd edition, Wiley, 2002. (p. 515, 584)

or CYLINDRICAL

③

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Microscopic species mass balance, constant thermal conductivity; cylindrical coordinates

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(4)

$$\left(\frac{1}{M_A}\right) \rho \frac{\partial w_A}{\partial t} = \rho D_{AB} \frac{\partial^2 w_A}{\partial z^2} \left(\frac{1}{M_A}\right)$$

$$\frac{\partial c_A}{\partial t} = D_{AB} \frac{\partial^2 c_A}{\partial z^2}$$

multiply by $\frac{1}{M_A}$ (both sides)

IC: $t = 0 \quad c_A = c_{A0} \quad \forall z$

BC: $z = 0 \quad c_A = c_{As} \quad \forall t > 0$

$z = \infty \quad c_A = c_{A0} \quad \forall t > 0$



O₂ Diffusion Solution:

The oxygen concentration as a function of time and depth into the water is given by:

$$\frac{c_{As} - c_A}{c_{As} - c_{A0}} = \text{erfc} \left(\frac{z}{2\sqrt{D_{AB}t}} \right) = \text{erfc} \zeta$$

$$\zeta \equiv \frac{z}{2\sqrt{D_{AB}t}}$$

Solution is from WRF (or text) or the literature.

Unsteady State Mass Transport

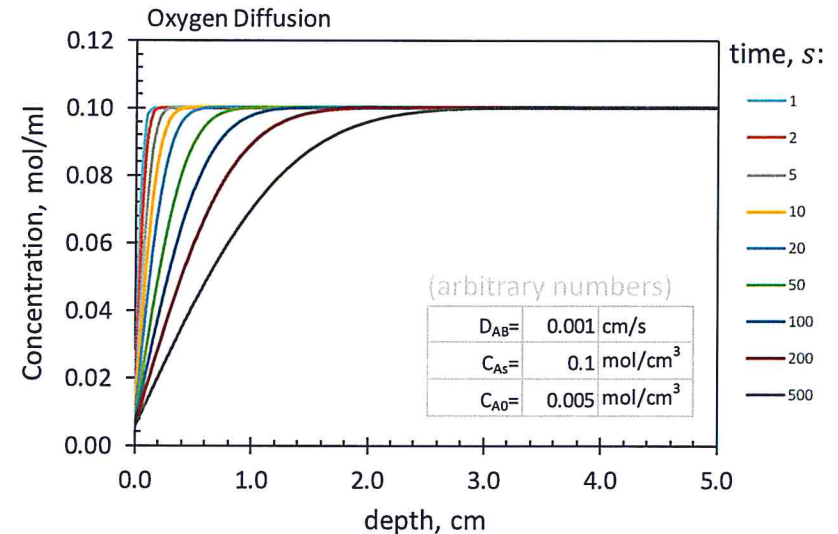
Unsteady State Diffusion in a Semi-Infinite Slab

$$\frac{\partial c_A}{\partial t} = D_{AB} \left(\frac{\partial^2 c_A}{\partial z^2} \right)$$

Initial condition: $t = 0 \quad c_A = c_{A0} \quad \forall z$

Boundary conditions:

$$x = 0 \quad c_A = c_{As} \quad t > 0$$

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$$\zeta \equiv \frac{z}{2\sqrt{D_{AB}t}}$$

This solution was a resource in the Danckwerts model for mass transfer; the short penetration time meant that the diffusion direction looked “infinite.”

Unsteady State Mass Transport

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Initial condition: $t = 0 \quad c_A = c_{A0} \quad \forall z$

Boundary conditions:

$$x = 0 \quad c_A = c_{As} \quad t > 0$$

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Mass Transport “Laws”

Another physical picture associated with penetration theory is “surface renewal” (Danckwerts)

- Turbulent flow
- Diffusing species only penetrates a short distance
- Due to chem rxn or short time of contact, t_{exp}
- Model as unsteady state molecular transport
- Danckwerts: bulk motion brings fresh liquid eddies from interior to the surface
- At the surface A is transferred as though B were stagnant and infinitely deep
- Works for falling film

$$k_c = \sqrt{\frac{4D_{AB}}{\pi t_{exp}}}$$
