




## CM3120 Transport/Unit Operations 2







**Professor Faith A. Morrison**  
 Department of Chemical Engineering  
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[www.chem.mtu.edu/~fmorriso/cm3120/cm3120.html](http://www.chem.mtu.edu/~fmorriso/cm3120/cm3120.html)


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## Where to start?

We seek to study **unsteady state heat transfer**.  
 Let's start by looking over several subjects that form the foundation for what we hope to study.

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These basic concepts are familiar. Cycling back can deepen our understanding (and help us put new concepts in context)

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## CM3120 Transport/Unit Operations 2

### Energy Balance Review



*Professor Faith A. Morrison*

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### CM2110/CM2120 - Review

#### Macroscopic Energy Balances

The physics:

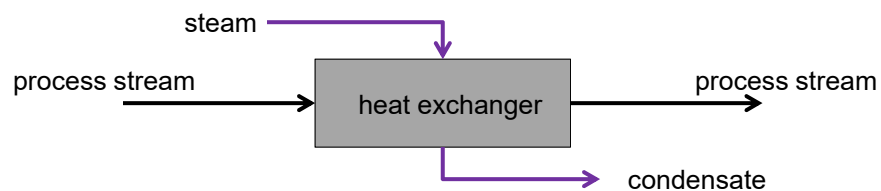
- Mass conservation
- Energy conservation

Open system energy balance

$$\Delta E_p + \Delta E_k + \Delta H = Q_{in} + W_{s,on} \quad (\text{out-in})$$

Closed system energy balance

$$\Delta E_p + \Delta E_k + \Delta U = Q_{in} + W_{on} \quad (\text{final-initial})$$



[www.chem.mtu.edu/~fmorriso/cm310/Energy\\_Balance\\_Notes\\_2008.pdf](http://www.chem.mtu.edu/~fmorriso/cm310/Energy_Balance_Notes_2008.pdf)

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**CM2110/CM2120 - Review**

**Review:**

*How do we decide what equations to use for what?*

- Closed system E-bal (first law of thermo)
- Open system E-bal ( $H = U + PV$ , flowing systems)
- Mechanical energy balance (SISO, steady, isothermal, no rxn, no phase change, little heat transferred)

Knowing what assumptions we are making means we understand our models. See handout for summary:

[www.chem.mtu.edu/~fmorriso/cm310/Energy\\_Balance\\_Notes\\_2008.pdf](http://www.chem.mtu.edu/~fmorriso/cm310/Energy_Balance_Notes_2008.pdf)

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**Macroscopic Energy Balances**

**Open system energy balance**  
 $\Delta E_p + \Delta E_k + \Delta H = Q_{in} + W_{s,on}$  (out-in)

**Closed system energy balance**  
 $\Delta E_p + \Delta E_k + \Delta U = Q_{in} + W_{on}$  (final-initial)

MEB:  $\frac{\Delta p}{\rho} + \frac{\Delta(v^2)}{2\alpha} + g\Delta z + F = \frac{W_{s,on}}{\dot{m}}$

**Notes:**

1.  $\Delta$  has different meanings in the three e-balances
2. Open-system:  

$$\Delta H = \sum_{outs} m_i \hat{h}_i - \sum_{ins} m_i \hat{h}_i$$
3. Use MEB if you can (easy); but not if it does not apply!

What **physics** determines how rapidly the heat transfers from the outside stream to the inside stream?

**Fourier's Law of Heat Conduction**

$$\frac{q_x}{A} = -k \frac{dT}{dx}$$

(for a homogeneous phase)

$\frac{q_x}{A}$  – heat flux=energy/area time)  
 $k$  – thermal conductivity  
 $\frac{dT}{dx}$  –temperature gradient

(the driving physics is **Brownian motion**: energy transports down  $\nabla T$  due to Brownian motion)

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**CM2110/CM2120 - Review**

**Macroscopic Energy Balances**

**Open system energy balance**  
 $\Delta E_p + \Delta E_k + \Delta H = Q_{in} + W_{s,on}$  (out-in)

**Closed system energy balance**  
 $\Delta E_p + \Delta E_k + \Delta U = Q_{in} + W_{on}$  (final-initial)

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**Heat Transfer Rate law: Fourier's law of Heat Conduction:**

makes reference to a coordinate system

$$\frac{q_x}{A} = -k \frac{dT}{dx}$$

Allows you to solve for temperature profiles

Gibbs notation:  $\frac{q}{A} = -k \nabla T$

Fourier's law  
3D

$$\vec{q} = \frac{q}{A} = \begin{pmatrix} -k \frac{\partial T}{\partial x} \\ -k \frac{\partial T}{\partial y} \\ -k \frac{\partial T}{\partial z} \end{pmatrix}_{xyz}$$

- Heat flows **down** a temperature gradient
- Flux is proportional to magnitude of temperature gradient

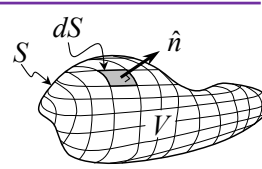
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To use the heat transport law (Fourier's law) we need to use the **microscopic** energy balance.

**Microscopic Energy Balance:**  
Equation of Thermal Energy



Microscopic **energy** balance written on an arbitrarily shaped volume, V, enclosed by a surface, S

Gibbs notation:  $\rho \left( \frac{\partial \hat{E}}{\partial t} + \underline{v} \cdot \nabla \hat{E} \right) = -\nabla \cdot \vec{q} + S_e$  general conduction

Gibbs notation:  $\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S_e$  Only Fourier conduction

(incompressible fluid, constant pressure, neglect  $\hat{E}_k, \hat{E}_p$ , viscous dissipation)

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[http://pages.mtu.edu/~fmorriso/cm310/energy\\_equation.html](http://pages.mtu.edu/~fmorriso/cm310/energy_equation.html)

**CM3110 - Review**

**CM3110  
REVIEW**

**Equation of Energy**  
(microscopic energy balance)

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S_e$$

rate of change

convection

source  
(energy generated per unit volume per time)

conduction  
(all directions)

velocity must satisfy equation of motion, equation of continuity

see handout for component notation

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The **Equation of Energy** in Cartesian, cylindrical, and spherical coordinates for Newtonian fluids of constant density, with source term S. Source could be electrical energy due to current flow, chemical energy, etc. Two cases are presented: the general case where thermal conductivity may be a function of temperature (vector flux  $\hat{q} = q/A$  appears in the equations); and the more usual case, where thermal conductivity is constant.

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**Microscopic energy balance**, in terms of flux; Gibbs notation

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = -\nabla \cdot \hat{q} + S$$

**Microscopic energy balance**, in terms of flux; Cartesian coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = - \left( \frac{\partial q_x}{\partial x} + \frac{\partial q_y}{\partial y} + \frac{\partial q_z}{\partial z} \right) + S$$

**Microscopic energy balance**, in terms of flux; cylindrical coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = - \left( \frac{1}{r} \frac{\partial (r \hat{q}_r)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial \hat{q}_\theta}{\partial \theta} + \frac{\partial \hat{q}_z}{\partial z} \right) + S$$

**Microscopic energy balance**, in terms of flux; spherical coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial T}{\partial \phi} \right) = - \left( \frac{1}{r^2} \frac{\partial (r^2 \hat{q}_r)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial (\hat{q}_\theta \sin \theta)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial \hat{q}_\phi}{\partial \phi} \right) + S$$

**Fourier's law of heat conduction**, Gibbs notation:  $\hat{q} = -k \nabla T$

**Fourier's law of heat conduction**, Cartesian coordinates:

$$\begin{pmatrix} \hat{q}_x \\ \hat{q}_y \\ \hat{q}_z \end{pmatrix}_{xyz} = \begin{pmatrix} -k \frac{\partial T}{\partial x} \\ -k \frac{\partial T}{\partial y} \\ -k \frac{\partial T}{\partial z} \end{pmatrix}_{xyz}$$

**Fourier's law of heat conduction**, cylindrical coordinates:

$$\begin{pmatrix} \hat{q}_r \\ \hat{q}_\theta \\ \hat{q}_z \end{pmatrix}_{r\theta z} = \begin{pmatrix} -k \frac{\partial T}{\partial r} \\ \frac{v}{r} \frac{\partial T}{\partial \theta} \\ -k \frac{\partial T}{\partial z} \end{pmatrix}_{r\theta z}$$

**Fourier's law of heat conduction**, spherical coordinates:

$$\begin{pmatrix} \hat{q}_r \\ \hat{q}_\theta \\ \hat{q}_\phi \end{pmatrix}_{r\theta\phi} = \begin{pmatrix} -k \frac{\partial T}{\partial r} \\ -k \frac{\partial T}{\partial \theta} \\ -\frac{k}{r \sin \theta} \frac{\partial T}{\partial \phi} \end{pmatrix}_{r\theta\phi}$$

The **Equation of Energy** for systems with **constant k**

**Microscopic energy balance**, constant thermal conductivity; Gibbs notation

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S$$

**Microscopic energy balance**, constant thermal conductivity; Cartesian coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; cylindrical coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; spherical coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial T}{\partial \phi} \right) = k \left( \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial T}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 T}{\partial \phi^2} \right) + S$$

Reference: F. A. Morrison, "Web Appendix to An Introduction to Fluid Mechanics," Cambridge University Press, New York, 2013. On the web at [www.chem.mtu.edu/~fmorriso/IFM\\_WebAppendixCD2013.pdf](http://www.chem.mtu.edu/~fmorriso/IFM_WebAppendixCD2013.pdf)

[http://pages.mtu.edu/~fmorriso/cm310/energy\\_equation.html](http://pages.mtu.edu/~fmorriso/cm310/energy_equation.html)
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**The Equation of Energy** in Cartesian, cylindrical, and spherical coordinates for Newtonian fluids of constant density, with source term  $S$ . Source could be electrical energy due to current flow, chemical energy, etc. Two cases are presented: the general case where thermal conductivity may be a function of temperature (vector flux  $\underline{\tilde{q}} = \underline{q}/A$  appears in the equations); and the more usual case, where thermal conductivity is constant.

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**Microscopic energy balance**, in terms of flux; Gibbs notation

$$\rho c_p \left( \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = -\nabla \cdot \underline{\tilde{q}} + S$$

**Microscopic energy balance**, in terms of flux; Cartesian coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = - \left( \frac{\partial \tilde{q}_x}{\partial x} + \frac{\partial \tilde{q}_y}{\partial y} + \frac{\partial \tilde{q}_z}{\partial z} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; Cartesian coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; cylindrical coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; spherical coordinates

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**The Equation of Energy** for systems with constant  $k$

**Microscopic energy balance**, constant thermal conductivity; Gibbs notation

$$\rho c_p \left( \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S$$

**Microscopic energy balance**, constant thermal conductivity; Cartesian coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; cylindrical coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; spherical coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial T}{\partial \phi} \right) = k \left( \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial T}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 T}{\partial \phi^2} \right) + S$$

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**Front side:**

- Micro E-bal in terms of flux  $\underline{\tilde{q}} \equiv \frac{\underline{q}}{A}$
- Fourier's law,  $\underline{\tilde{q}} = -k \nabla T$

Fourier's law of heat conduction, cylindrical coordinates:  $\left( \begin{matrix} \tilde{q}_r \\ \tilde{q}_\theta \\ \tilde{q}_z \end{matrix} \right) = \left( \begin{matrix} -k \frac{\partial T}{\partial r} \\ -k \frac{1}{r} \frac{\partial T}{\partial \theta} \\ -k \frac{\partial T}{\partial z} \end{matrix} \right)$

Fourier's law of heat conduction, spherical coordinates:  $\left( \begin{matrix} \tilde{q}_r \\ \tilde{q}_\theta \\ \tilde{q}_\phi \end{matrix} \right) = \left( \begin{matrix} -k \frac{\partial T}{\partial r} \\ -k \frac{1}{r} \frac{\partial T}{\partial \theta} \\ -k \frac{1}{r \sin \theta} \frac{\partial T}{\partial \phi} \end{matrix} \right)$

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**Microscopic energy balance**, in terms of flux; Gibbs notation

$$\rho c_p \left( \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = -\nabla \cdot \underline{\tilde{q}} + S$$

**Microscopic energy balance**, in terms of flux; Cartesian coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = - \left( \frac{\partial \tilde{q}_x}{\partial x} + \frac{\partial \tilde{q}_y}{\partial y} + \frac{\partial \tilde{q}_z}{\partial z} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; Cartesian coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; cylindrical coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

**Microscopic energy balance**, constant thermal conductivity; spherical coordinates

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**The Equation of Energy** for systems with constant  $k$

**Microscopic energy balance**, constant thermal conductivity; Gibbs notation

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**Microscopic energy balance**, constant thermal conductivity; Cartesian coordinates

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**Microscopic energy balance**, constant thermal conductivity; cylindrical coordinates

$$\rho c_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

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**Back side:**

- Micro E-bal in terms of temperature (Fourier's law incorporated)

Fourier's law of heat conduction, cylindrical coordinates:  $\left( \begin{matrix} \tilde{q}_r \\ \tilde{q}_\theta \\ \tilde{q}_z \end{matrix} \right) = \left( \begin{matrix} -k \frac{\partial T}{\partial r} \\ -k \frac{1}{r} \frac{\partial T}{\partial \theta} \\ -k \frac{\partial T}{\partial z} \end{matrix} \right)$

Fourier's law of heat conduction, spherical coordinates:  $\left( \begin{matrix} \tilde{q}_r \\ \tilde{q}_\theta \\ \tilde{q}_\phi \end{matrix} \right) = \left( \begin{matrix} -k \frac{\partial T}{\partial r} \\ -k \frac{1}{r} \frac{\partial T}{\partial \theta} \\ -k \frac{1}{r \sin \theta} \frac{\partial T}{\partial \phi} \end{matrix} \right)$

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## Microscopic Energy Balance

CM3110  
REVIEWThe **Equation of Energy** for systems with **constant  $k$** 

Microscopic energy balance, constant thermal conductivity; Gibbs notation

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S$$

Microscopic energy balance, constant thermal conductivity; Cartesian coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

Microscopic energy balance, constant thermal conductivity; cylindrical coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + v_z \frac{\partial T}{\partial z} \right) = k \left( \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right) + S$$

Microscopic energy balance, constant thermal conductivity; spherical coordinates

$$\rho \hat{c}_p \left( \frac{\partial T}{\partial t} + v_r \frac{\partial T}{\partial r} + \frac{v_\theta}{r} \frac{\partial T}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial T}{\partial \phi} \right) = k \left( \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial T}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 T}{\partial \phi^2} \right) + S$$

[http://pages.mtu.edu/~fmorriso/cm310/energy\\_equation.html](http://pages.mtu.edu/~fmorriso/cm310/energy_equation.html)

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## Fourier's Law of Heat Conduction

CM3110  
REVIEWFourier's law of heat conduction, Gibbs notation:  $\tilde{q} = -k \nabla T$ 

Fourier's law of heat conduction, Cartesian coordinates:

$$\begin{pmatrix} \tilde{q}_x \\ \tilde{q}_y \\ \tilde{q}_z \end{pmatrix}_{xyz} = \begin{pmatrix} -k \frac{\partial T}{\partial x} \\ -k \frac{\partial T}{\partial y} \\ -k \frac{\partial T}{\partial z} \end{pmatrix}_{xyz}$$

$$\tilde{q}_x = \frac{q_x}{A}$$

Fourier's law of heat conduction, cylindrical coordinates:

$$\begin{pmatrix} \tilde{q}_r \\ \tilde{q}_\theta \\ \tilde{q}_z \end{pmatrix}_{xyz} = \begin{pmatrix} -k \frac{\partial T}{\partial r} \\ -\frac{k}{r} \frac{\partial T}{\partial \theta} \\ -k \frac{\partial T}{\partial z} \end{pmatrix}_{r\theta z}$$

Fourier's law of heat conduction, spherical coordinates:

$$\begin{pmatrix} \tilde{q}_r \\ \tilde{q}_\theta \\ \tilde{q}_\phi \end{pmatrix}_{xyz} = \begin{pmatrix} -k \frac{\partial T}{\partial r} \\ -\frac{k}{r} \frac{\partial T}{\partial \theta} \\ -\frac{k}{r \sin \theta} \frac{\partial T}{\partial \phi} \end{pmatrix}_{r\theta\phi}$$

[http://pages.mtu.edu/~fmorriso/cm310/energy\\_equation.html](http://pages.mtu.edu/~fmorriso/cm310/energy_equation.html)

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We will need **boundary conditions** on temperature

**CM3110 REVIEW**

**Example 1: Heat flux in a rectangular solid – Temperature BC**

*What is the steady state temperature profile in a rectangular slab if one side is held at  $T_1$  and the other side is held at  $T_2$ ?*

**Assumptions:**

- wide, tall slab
- steady state

$\frac{q_x}{A}$

HOT SIDE

COLD SIDE

$T_1 > T_2$

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We will need **boundary conditions** on temperature

**CM3110 REVIEW**

**Example 1: Heat flux in a rectangular solid – Temperature BC**

*What is the steady state temperature profile in a rectangular slab if one side is held at  $T_1$  and the other side is held at  $T_2$ ?*

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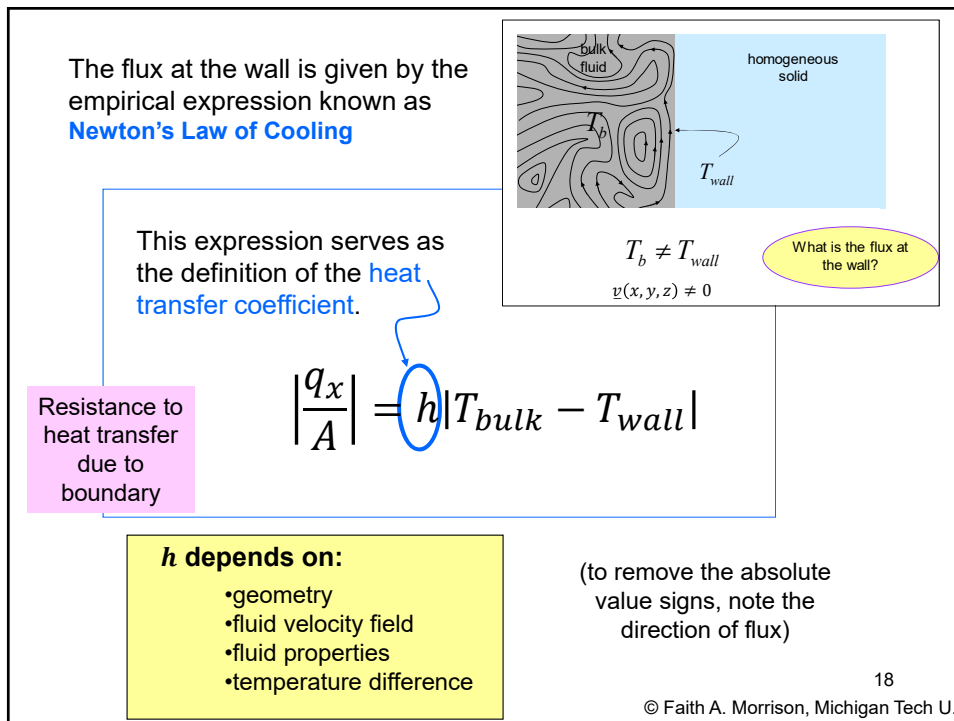
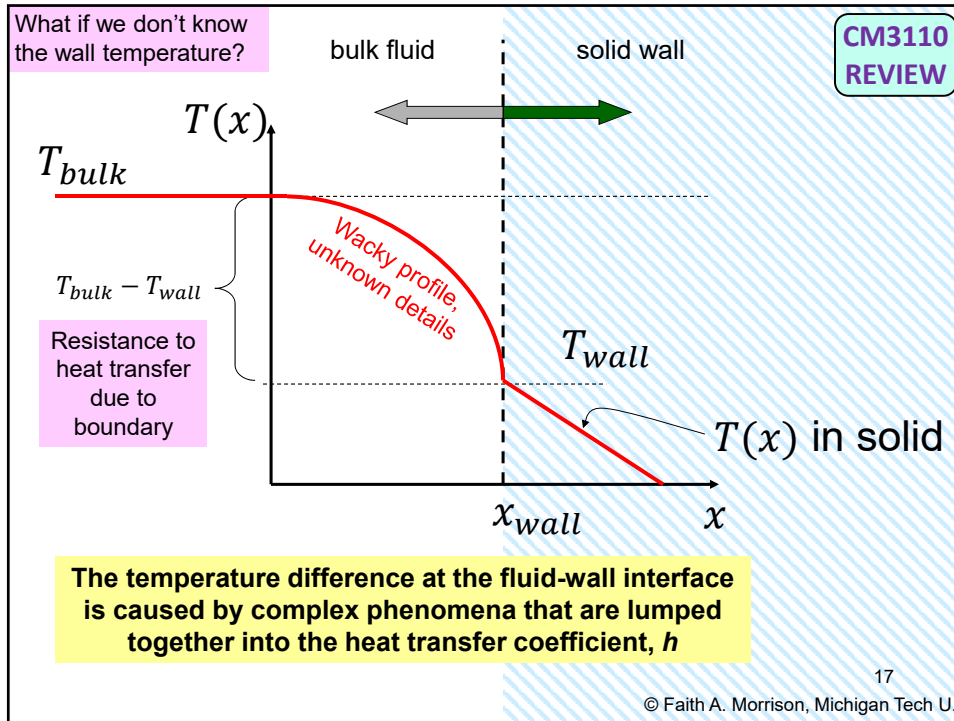
HOT SIDE

What if we don't know the wall temperature?

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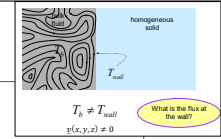




**Note:**

Large values of  $h$  are associated with small  $|T_{bulk} - T_{wall}|$

The flux at the wall is given by the empirical expression known as **Newton's Law of Cooling**



This expression serves as the definition of the **heat transfer coefficient**.

Resistance to heat transfer due to boundary

$$\left| \frac{q_x}{A} \right| = h |T_{bulk} - T_{wall}|$$

**h depends on:**

- geometry
- fluid velocity field
- fluid properties
- temperature difference

(to remove the absolute value signs, note the direction of flux)

$$\left| \frac{q_x}{A} \right| = h |T_{bulk} - T_{wall}|$$

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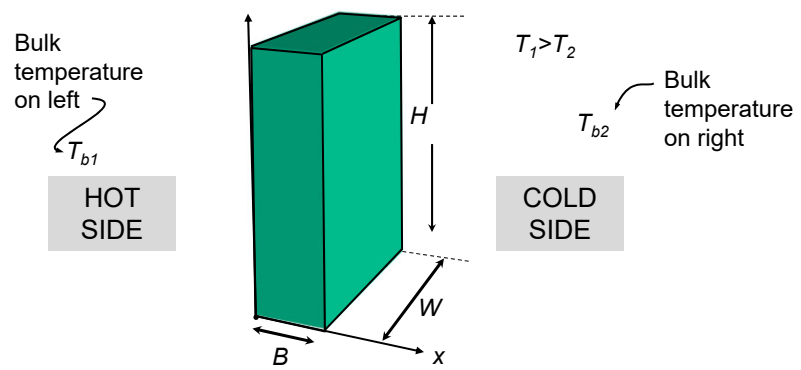
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Newton's law of cooling BC

CM3110  
REVIEW

**Example 2: Heat flux in a rectangular solid – Newton's law of cooling BC**

*What is the steady state temperature profile in a rectangular slab if the fluid on one side is held at  $T_{b1}$  and the fluid on the other side is held at  $T_{b2}$ ?*



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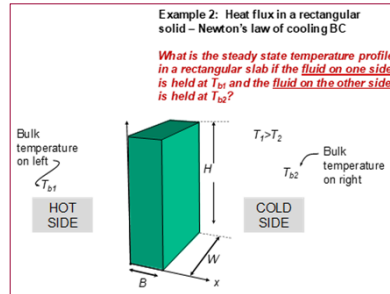
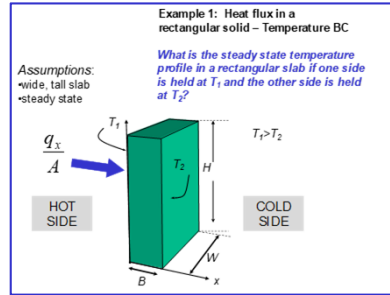
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**We will need boundary conditions on temperature**

- Commonly,
  - Temperature boundary conditions
 
$$x = 0 \quad T = T_1$$

$$x = B \quad T = T_2$$
  - Newton's law of cooling (flux) boundary conditions
 
$$\left. \frac{q_x}{A} \right|_{x=0} = h_1(T_{b1} - T(0))$$

$$\left. \frac{q_x}{A} \right|_{x=B} = h_1(T(B) - T_{b2})$$
- Sometimes, we know a value for the flux at the wall.



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**Let's carry out an example of 1D, steady heat transfer**

**CM3110 REVIEW**

CM3110  
Transport I  
Part II: Heat Transfer

**MichiganTech**

**One-Dimensional Heat Transfer**  
(part 1: rectangular slab)

Simple problems that allow us to identify the physics

**Professor Faith Morrison**  
Department of Chemical Engineering  
Michigan Technological University

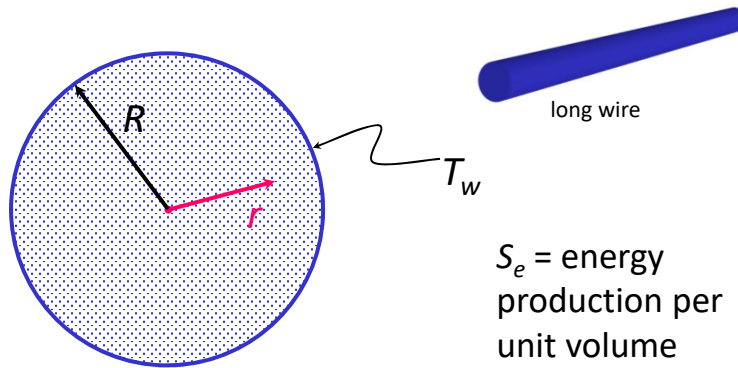
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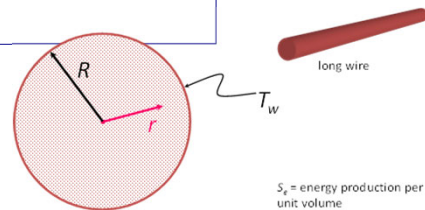
**Example 3: Heat Conduction with Generation**

What is the steady state temperature profile in a wire if heat is generated uniformly throughout the wire at a rate of  $S_e$  W/m<sup>3</sup> and the outer radius is held at  $T_w$ ? What is the flux?



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**Example: Heat conduction with generation**

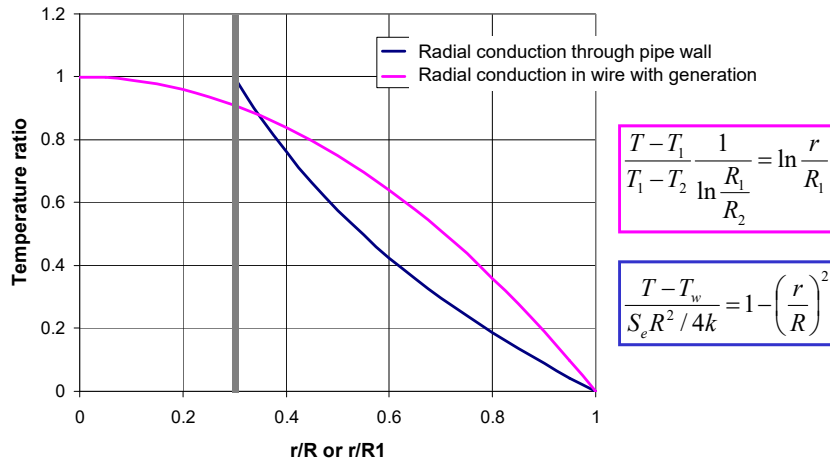


You try.

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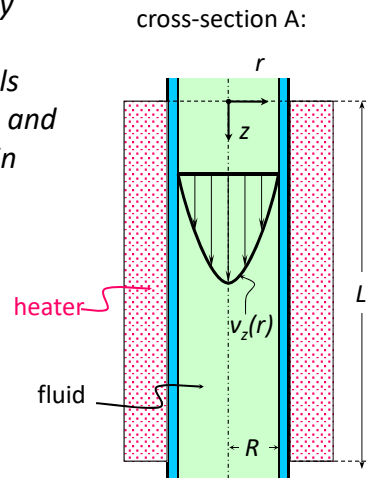
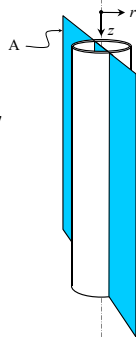
### Compare radial conduction solutions



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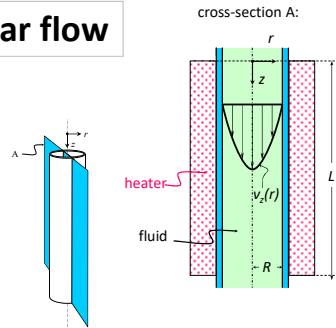
**Example 4: Wall heating of laminar flow.** What is the steady state temperature profile in a flowing fluid in a tube if the walls are heated (constant flux,  $q_1/A$ ) and if the fluid is a Newtonian fluid in laminar flow?

**assume:**  
constant viscosity



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**Example 4: Wall heating of laminar flow**

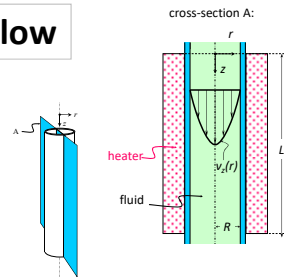


You try.

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**Example 4: Wall heating of laminar flow**



We need to solve this partial differential equation:

$$\frac{1}{r} \frac{\partial}{\partial r} \left( \frac{\partial T}{\partial r} r \right) + \frac{\partial}{\partial z} \left( \frac{\partial T}{\partial z} \right) - \frac{\rho \hat{c}_p}{k} v_z(r) \frac{\partial T}{\partial z} = 0$$

with

$$v_z = \frac{\Delta p}{4\mu L} (R^2 - r^2)$$

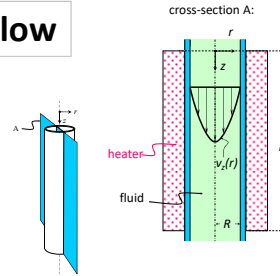
and with the appropriate boundary conditions. To see the solution go to:

- R. Siegel, E. M. Sparrow, T. M. Hallman, *Appl. Science Research* A7, 386-392 (1958)
- R. B. Bird, W. Stewart, and E. Lightfoot (BSL), *Transport Phenomena*, Wiley, 1960, p295.

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**Example 4: Wall heating of laminar flow**



We need to solve this partial differential equation:

$$\frac{1}{r} \frac{\partial}{\partial r} \left( \frac{\partial T}{\partial r} r \right) + \frac{\partial}{\partial z} \left( \frac{\partial T}{\partial z} \right) - \frac{\rho \hat{c}_p}{k} v_z(r) \frac{\partial T}{\partial z} = 0$$

with

$$v_z = \frac{\Delta p}{4\mu L} (R^2 - r^2)$$

and with the appropriate boundary conditions the solution goes to:

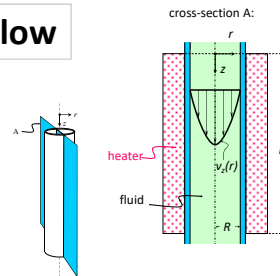
**What are the BCs?**

- R. Siegel, E. M. Sparrow, *ASME JOURNAL OF HEAT TRANSFER*, Vol. 77, No. 4, pp. 271-278, 1955
- R. B. Bird, W. Stewart, and E. Lightfoot (BSL), *Transport Phenomena*, Wiley, 1960, p295.

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**Example 4: Wall heating of laminar flow**



Answer:  
Neglect axial conduction

BC:

$$\begin{aligned} r = 0 & \quad T = \text{finite} \\ r = R & \quad \frac{q_r}{A} = \frac{q_1}{A} \\ z = 0 & \quad T = T_0 \end{aligned}$$

See BSL p295

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## 1D, steady heat transfer

CM3110  
REVIEW

### SUMMARY

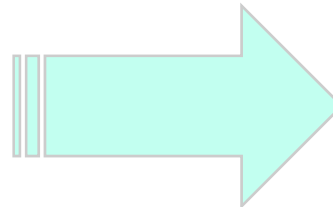
- Microscopic energy balance
- Transport law
- Newton's law of cooling (fluid boundary)
- Know the assumptions that simplify the model of the problem
- Solve with appropriate boundary conditions
- Check that assumptions are valid when using the solution

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## NEXT: Unsteady State Heat Transfer

The thumbnail shows a presentation slide with a blue header "CM3120 Transport/Unit Operations 2". Below the header is a red oval containing the text "Unsteady State Heat Transfer". To the right of the oval is a photograph of a laboratory setup with various pipes and tanks. Below the photograph is a small portrait of Professor Faith A. Morrison, with her name and affiliation: "Professor Faith A. Morrison, Department of Chemical Engineering, Michigan Technological University". At the bottom left of the slide is the URL "www.chem.mtu.edu/~fmorriso/cm3120/cm3120.html" and at the bottom right is the copyright notice "© Faith A. Morrison, Michigan Tech U."



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