

Given that $\int_0^1 (f-f^3) d\eta$ is some constant, α

We restate the integral b.l. eqn as

$$\rho U^2 \alpha \frac{d\delta}{dx} = \tau_0$$

We may rewrite

$$\tau_0 = \mu \left[\frac{\partial v_x}{\partial y} \right]_{y=0} = \frac{\mu U}{\delta} \left(\frac{df}{d\eta} \right)_{\eta=0} = \frac{\mu U}{\delta(x)} \beta$$

where $\beta \equiv \left(\frac{\partial f}{\partial \eta} \right)_{\eta=0}$ is some constant also.

Thus, we now have a governing eqn. for $\delta(x)$.

$$\rho U^2 \alpha \frac{d\delta}{dx} = \frac{\mu U}{\delta} \beta$$

$$\delta \frac{d\delta}{dx} = \frac{\mu \beta}{\rho U \alpha}$$

$$x=0, \delta=0.$$

Soln is $\boxed{\delta = \left(\frac{2\mu\beta}{\rho U \alpha} x \right)^{1/2}}$ thus δ varies as \sqrt{x} !

An approximation to $v_x(\frac{y}{\delta(x)})$ is.

$$\frac{v_x}{u} = \frac{3}{2} \frac{y}{\delta(x)} - \frac{1}{2} \left(\frac{y}{\delta(x)} \right)^3 \rightarrow f = \frac{3}{2} \eta - \frac{1}{2} \eta^3$$

this solution agrees with the boundary conditions.

$$\text{BC1 } y=0, \quad v_x=0$$

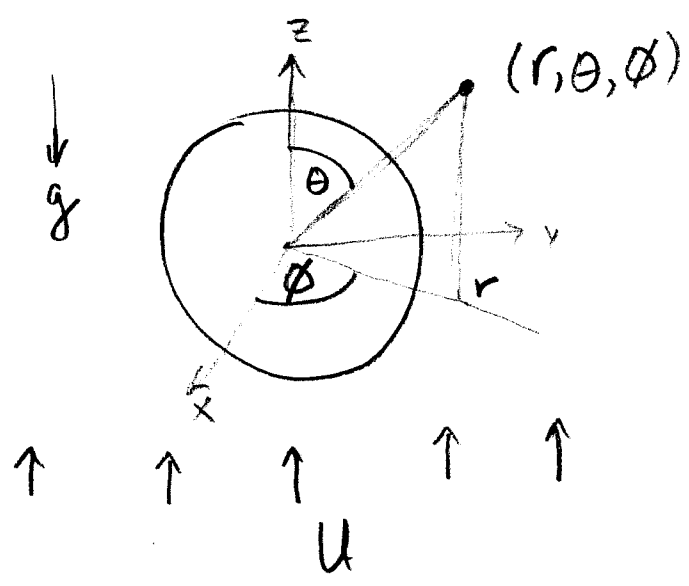
$$\text{BC2 } y=\delta(x), \quad \frac{\partial v_x}{\partial y} = 0 \quad \text{"smoothness" condition}$$

$$\text{BC3 } y=\delta(x), \quad v_x=u$$

Calc. α and β and then solve for $\delta(x)$.

Then Calc. F_x , the Force on the Plate!

Slow Flow Around a Solid Sphere: 2-dimensional Flow: Stokes' Law Derivation (Middleman, 168-172)



We "fix" the r, θ, ϕ coordinate system to the falling sphere. U appears to the sphere as a uniform constant velocity field at steady state (when drag force = gravity force).

Symmetry dictates that $v_\phi = 0$ Mon 2/6/06

But we expect $v_\theta = v_\theta(r, \theta)$, $v_r = v_r(r, \theta)$

For a Newtonian fluid of constant ρ we may write

Continuity Eqn: Table B.4.

$$\frac{\partial \rho}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} (\rho r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\rho v_\theta \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} (\rho v_\phi) = 0$$

ρ
Const.

$$\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (v_\theta \sin \theta) = 0$$