



CM3110
Transport I
Part II: Heat Transfer



Michigan Tech



Radiation Heat Transfer


- In Unit Operations
- Heat Shields

Professor Faith Morrison

Department of Chemical Engineering
Michigan Technological University

© Faith A. Morrison, Michigan Tech U.

CM3110
Transport Processes and Unit Operations I
Part 2: Heat Transfer



Michigan Tech

Summary *(Part 2 thus far)*

Within homogeneous phases:

- Microscopic Energy Balances
- 1D Steady solutions

rectangular: $\frac{q_x}{A} = C_1$
 $T = ax + b$

cylindrical: $\frac{q_r}{A} = \frac{C_1}{r}$
 $T = a \ln x + b$

$$\rho \hat{c}_p \left(\frac{\partial T}{\partial t} + \mathbf{v} \cdot \nabla T \right) = k \nabla^2 T + S_e$$


conduction

- Temperature and *Newton's law of cooling* boundary conditions (if h is supplied)
- Unsteady solutions (from literature)
 - ✓ Carslaw and Jeager
 - ✓ Heisler charts

2

© Faith A. Morrison, Michigan Tech U.

CM3110
Transport Processes and Unit Operations I
Part 2: Heat Transfer



Michigan Tech

Summary *(Part 2 thus far)*

Across phase boundaries:

- Microscopic Energy, Momentum, and Mass Balances

Micro momentum:

$$\rho \left(\frac{\partial \underline{v}}{\partial t} + \underline{v} \cdot \nabla \underline{v} \right) = -\nabla p + \mu \nabla^2 \underline{v} + \rho \underline{g}$$

Micro energy:

$$\rho \hat{c}_p \left(\frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S_e$$

convection

- Simultaneous effects (complex)
- Solutions are difficult to obtain (and often not really necessary)
 ⇒ use **dimensional analysis** and expts to obtain **h**
- *h* Data correlations for:
 - ✓ forced convection
 - ✓ natural convection
 - ✓ evaporation/condensation

→
radiation

}
phase change

One more type of heat transfer

© Faith A. Morrison, Michigan Tech U. ³

Radiation versus Conduction and Convection

Continuum view

- Conduction is caused by macroscopic temperature gradients
- Convection is caused by macroscopic flow
- Radiation? **NO CONTINUUM EXPLANATION**

© Faith A. Morrison, Michigan Tech U. ⁴

Continuum versus Molecular description of matter

A continuum is infinitely divisible

Real matter is not a continuum; at small enough length scales, molecules are discrete.

5
© Faith A. Morrison, Michigan Tech U.

Radiation versus Conduction and Convection

Continuum view

- Conduction is caused by macroscopic temperature gradients
- Convection is caused by macroscopic flow
- Radiation? **NO CONTINUUM EXPLANATION**

Molecular view

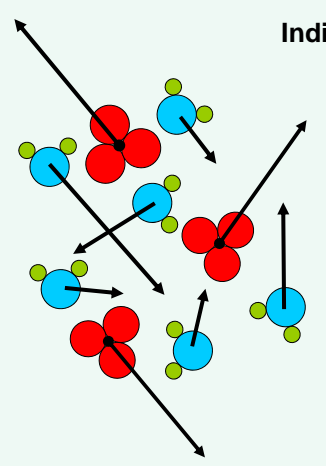
- Conduction—Brownian motion
- Convection—flow
- Radiation is caused by changes in electron energy states in molecules and atoms

There is also, of course, a molecular explanation of these effects, since we know that matter is made of atoms and molecules

Molecular view

© Faith A. Morrison, Michigan Tech U.

Molecular view



Individual molecules carry:

- chemical identity
- macroscopic velocity (speed and direction)
- internal energy (Brownian velocity)

When they undergo **Brownian motion** within an inhomogeneous mixture, they cause:

- diffusion** (mass transport)
- exchange** of momentum (viscous transport)
- conduction** (energy transport)

7

© Faith A. Morrison, Michigan Tech U.

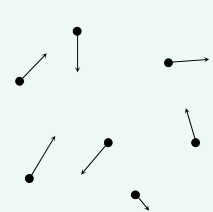
Molecular view

Kinetic Theory J. C. Maxwell, L. Boltzmann, 1860

- Molecules are in constant motion (Brownian motion)
- Temperature is related to $E_{k, av}$ of the molecules

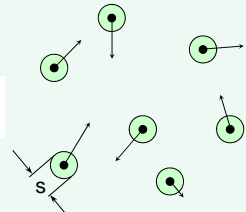
Simplest model

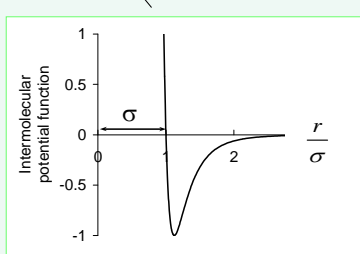
- no particle volume
- no intermolecular forces



More realistic model

- finite particle volume
- intermolecular forces





Intermolecular potential function vs r/σ . The graph shows a vertical asymptote at $r/\sigma = 1$, a minimum at $r/\sigma = 2$, and a horizontal asymptote at 0 for $r/\sigma > 2$.

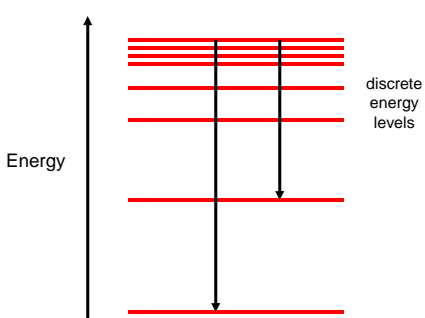
8

© Faith A. Morrison, Michigan Tech U.

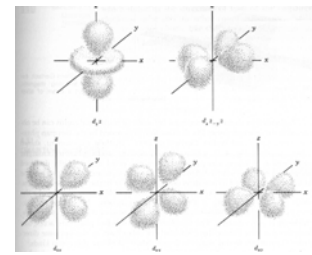
Molecular view

But, there is more to molecular energy than just Brownian motion...

- In atoms and molecules, electrons can exist in multiple, discrete energy states
- Transfers between energy states are accompanied by an emission of radiation



discrete energy levels



Sienko and Plane, Chemistry: Principles and Applications, McGraw Hill, 1979

Quantum Mechanics

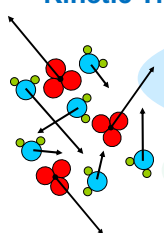
Radiation Heat Transfer is related to these non-Brownian mechanisms.

9

© Faith A. Morrison, Michigan Tech U.

Molecular view

Kinetic Theory

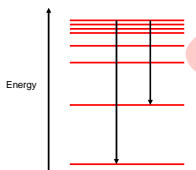


Is based on **Brownian motion** (molecules in constant motion proportional to their temperature)

Predicts that properties that are carried by individual molecules (chemical identity, momentum, average kinetic energy) will be transported **DOWN** gradients in these quantities.

⇒ Gradient transport laws are due to Brownian motion

Heat Transfer by Radiation



discrete energy levels

Is due to the release of energy stored in molecules that is **NOT** related to average kinetic energy (temperature), but rather to changing populations of excited states.

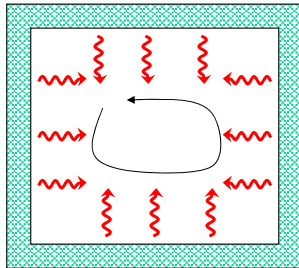
⇒ Radiation is **NOT** a Brownian effect

10

© Faith A. Morrison, Michigan Tech U.

How does this relate to chemical engineering?

Consider a furnace with an internal blower:



There is heat transfer due to convection:

$$q_{convection} = hA(T_s - T_b)$$

(Use correlations)

surface temp

Bulk temp

There is also heat transfer due to radiation.

$$q_{total} = q_{conv} + q_{rad}$$

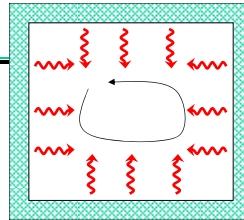
Where do we get q_{rad} ?

11

© Faith A. Morrison, Michigan Tech U.

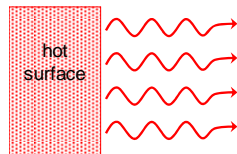
Where do we get $q_{radiation}$?

$$q_{total} = q_{conv} + q_{rad}$$

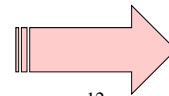


Answer:

Heat transfer due to radiation



We need to look into the physics of this mode of heat transfer.

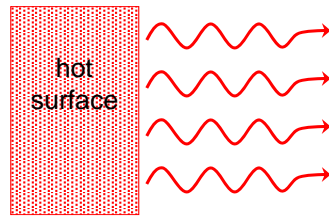


12

© Faith A. Morrison, Michigan Tech U.

Radiation

- does not require a medium to transfer energy (works in a vacuum)
- travels at the speed of light, $c = 3 \times 10^{10} \text{ cm/s}$
- travels as a wave; differs from x-rays, light, only by wavelength, λ
- radiation is important when temperatures are high



examples:

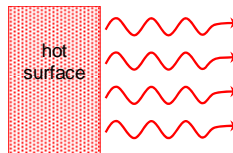
- the sun
- home radiator
- hot walls in vacuum oven
- heat exchanger walls when ΔT is high and a vapor film has formed

$$\frac{q}{A} \propto T^4$$

Note: **absolute temperature units**

13

© Faith A. Morrison, Michigan Tech U.



$$\frac{q}{A} \propto T^4$$

Why is radiation flux related to temperature and not to something else?

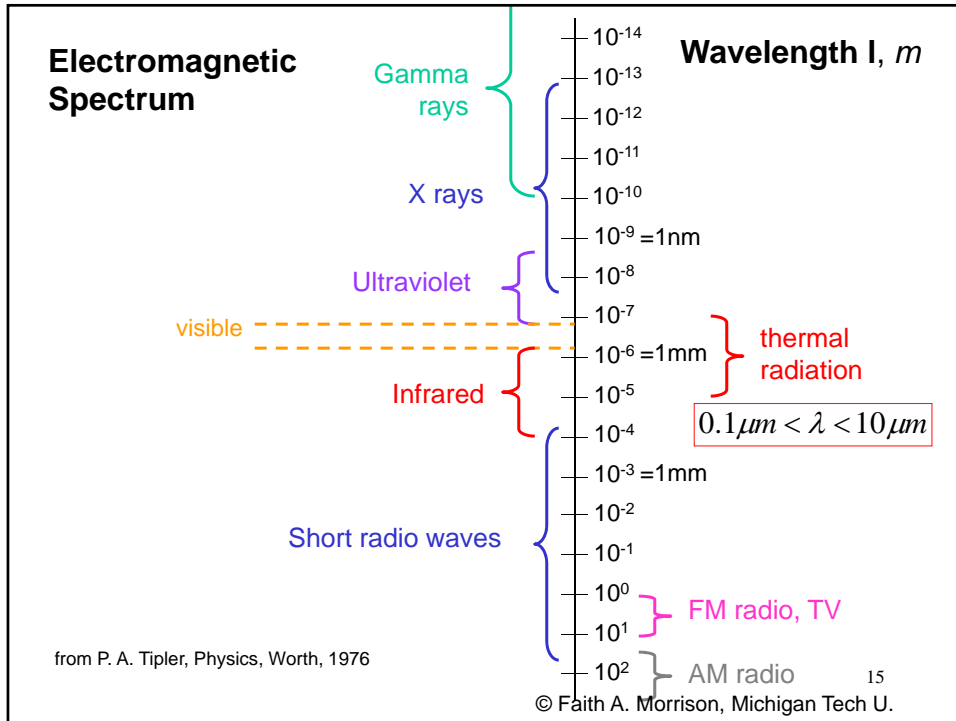
(From kinetic theory, temperature is related to average kinetic energy)

Answer:

- As a molecule gains energy, it **both** speeds up (increases average kinetic energy) and increases its population of excited states.
- The increase in **average kinetic energy** is reflected in temperature (directly proportional), and heat transfer through conduction.
- The increase in number of electrons in **excited states** is reflected in increased radiation heat flux. Electrons enter excited states in proportion to **absolute T^4** .

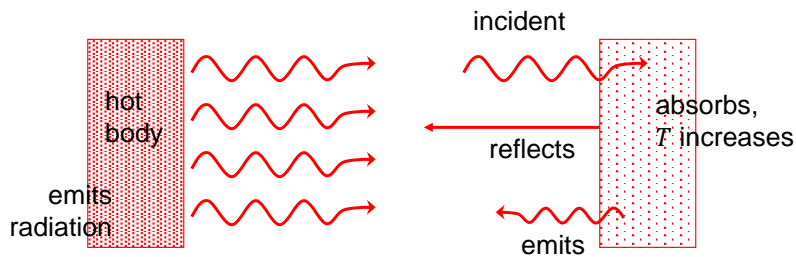
14

© Faith A. Morrison, Michigan Tech U.



What causes energy transfer by radiation?

- energy hits surface
- pushes some molecules into an excited state
- when the molecules/atoms relax from the excited state, they emit radiation



$$\frac{q_{emit}}{A} \propto T^4$$

$$\alpha = absorptivity$$

$$\alpha \equiv \frac{q_{absorbed}}{q_{incident}} < 1$$

Absorption

$\alpha = \text{absorptivity}$

$$\alpha \equiv \frac{q_{\text{absorbed}}}{q_{\text{incident}}} < 1$$

In general, absorptivity α is a function of wavelength

$\alpha = \alpha(\lambda)$

gray body: a body for which α is constant (does not depend on λ)
black body: a body for which $\alpha = 1$, i.e. absorbs all incident radiation

17
© Faith A. Morrison, Michigan Tech U.

Emission

$\varepsilon = \text{emissivity}$

$$\varepsilon \equiv \frac{q_{\text{emitted}}}{q_{\text{emitted, blackbody}}} < 1$$

gray body: a body for which α is constant
black body: a body for which $\alpha = 1$

$\alpha = \text{absorptivity}$

$$\alpha \equiv \frac{q_{\text{absorbed}}}{q_{\text{incident}}} < 1$$

Kirchhoff's Law: emissivity equals absorptivity at the same temperature

$\alpha = \varepsilon$

the fraction of energy absorbed by a material = the relative amount of energy emitted from that material compared to a black body

true for black and non-black solid surfaces

18
© Faith A. Morrison, Michigan Tech U.

$\epsilon = \text{emissivity}$
 $\epsilon \equiv \frac{q_{emitted}}{q_{emitted,black\ body}} < 1$

Black Bodies

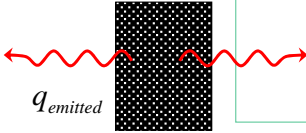
Stefan-Boltzmann Law: the amount of energy emitted by a black body is proportional to T^4

$$\frac{q_{emitted,black\ body}}{A} = \sigma T^4$$

NOTE: absolute temperature

$$\sigma = 0.1712 \times 10^{-8} \frac{BTU}{h\ ft^2R^4}$$

$$\sigma = 5.676 \times 10^{-8} \frac{W}{m^2K^4}$$



$q_{emitted}$

19

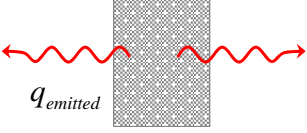
© Faith A. Morrison, Michigan Tech U.

Non-Black Bodies

$\epsilon = \text{emissivity}$
 $\epsilon \equiv \frac{q_{emitted}}{q_{emitted,black\ body}}$

$$\frac{q_{emitted,non-black\ body}}{A} = \epsilon q_{emitted,black\ body}$$

$$= \epsilon \sigma T^4$$



$q_{emitted}$

Stefan-Boltzmann:

$$\frac{q_{emitted,black\ body}}{A} = \sigma T^4$$

Energy emitted by a non-black body

$$\frac{q_{emitted,non-black\ body}}{A} = \epsilon \sigma T^4$$

20

© Faith A. Morrison, Michigan Tech U.

Radiation

Summary:

- Absorptivity, α
 - gray body: $\alpha = \text{constant}$
 - black body: $\alpha = 1$
- Emissivity, ε

$$q_{emit} = \varepsilon q_{emit,blackbody}$$
- Kirchoff's law: $\alpha = \varepsilon$
- Stefan-Boltzman law

$$\frac{q_{emit,blackbody}}{A} = \sigma T^4$$

$$\sigma = 0.1712 \times 10^{-8} \frac{BTU}{h ft^2 R^4}$$

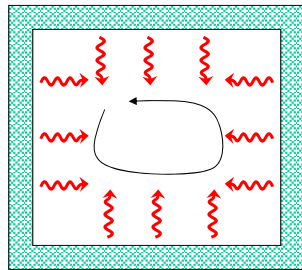
$$\sigma = 5.676 \times 10^{-8} \frac{W}{m^2 K^4}$$

21

© Faith A. Morrison, Michigan Tech U.

How does this relate to chemical engineering?

Consider a furnace with an internal blower:



There is heat transfer due to

convection:

$$q_{convection} = hA(T_s - T_b)$$

(Use correlations)

There is also heat transfer due to

radiation:

$$q_{radiation} = h_{rad}A(T_s - T_b)$$

$$q_{total} = q_{conv} + q_{rad}$$

$$q_{total} = (h_{conv} + h_{rad})A(T_s - T_b)$$

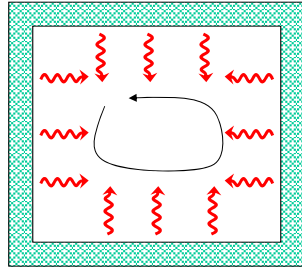
Because these two types of physics **do not interact**, we can just add the effects.

22

© Faith A. Morrison, Michigan Tech U.

How does this relate to chemical engineering?

Consider a furnace with an internal blower:



There is heat transfer due to convection:

$$q_{convection} = hA(T_s - T_b)$$

(Use correlations)

There is also heat transfer due to radiation:

$$q_{radiation} = h_{rad}A(T_s - T_b)$$

$$q_{total} = q_{conv} + q_{rad}$$

$$q_{total} = (h_{conv} + h_{rad})A(T_s - T_b)$$

Because these two types of physics **do not interact**, we can just add the effects.

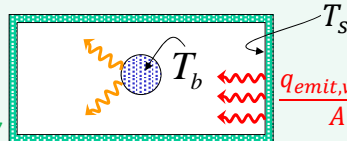
Where do we get h_{rad} ?

© Faith A. Morrison, Michigan Tech U.

23

Where do we get h_{rad} ?

$$\frac{q_{emit,body}}{A} = \epsilon\sigma T_{body}^4$$



$$\frac{q_{emit,walls}}{A} = \sigma T_s^4$$

black body, $\epsilon_{walls} \approx 1$

net flux to object:

$$\frac{q_{net}}{A} = \frac{q_{absorbed}}{A} - \frac{q_{emit}}{A}$$

Net flux to the body = Heat/area absorbed - Heat/area emitted

© Faith A. Morrison, Michigan Tech U.

24

Where do we get h_{rad} ?

$$\frac{q_{emit,body}}{A} = \epsilon \sigma T_{body}^4$$

$\frac{q_{emit,walls}}{A} = \sigma T_s^4$
black body, $\epsilon_{walls} \approx 1$

net flux to object:

$$\frac{q_{net}}{A} = \frac{q_{absorbed}}{A} - \frac{q_{emit}}{A}$$

$\frac{q_{absorbed,body}}{A} = \alpha \frac{q_{emit,walls}}{A}$

$\frac{q_{absorbed,body}}{A} = \alpha \sigma T_s^4 = \epsilon \Big|_{T_s} \sigma T_s^4$

$\alpha = \epsilon$, using Kirchhoff's law
 energy emitted by walls, which are acting as a black body

Heat/area absorbed

25

© Faith A. Morrison, Michigan Tech U.

Where do we get h_{rad} ?

$$\frac{q_{emit,body}}{A} = \epsilon \sigma T_{body}^4$$

$\frac{q_{emit,walls}}{A} = \sigma T_s^4$

net flux to object:

$$\frac{q_{net}}{A} = \frac{q_{absorbed}}{A} - \frac{q_{emit}}{A}$$

$$\frac{q_{emit,body}}{A} = \epsilon \Big|_{T_{body}} \sigma T_{body}^4$$

Object is a gray body,
 $\epsilon_{body} = \epsilon$

$$\frac{q_{emit,body}}{A} = \epsilon \Big|_{T_s} \sigma T_{body}^4$$

assuming: $\epsilon|_{T_s} \approx \epsilon|_{T_{body}}$

Heat/area emitted

26

© Faith A. Morrison, Michigan Tech U.

Where do we get h_{rad} ?

$$\frac{q_{emit,body}}{A} = \epsilon \sigma T_{body}^4$$

$$\frac{q_{emit,walls}}{A} = \sigma T_s^4$$

net flux to object:

$$\frac{q_{net}}{A} = \frac{q_{absorbed}}{A} - \frac{q_{emit}}{A}$$

Net flux to the body = Heat/area absorbed - Heat/area emitted

$$\frac{q_{net}}{A} = \epsilon \Big|_{T_s} (\sigma T_s^4 - \sigma T_{body}^4)$$

assuming: $\epsilon|_{T_s} \approx \epsilon|_{T_{body}}$

27

© Faith A. Morrison, Michigan Tech U.

Finally, calculate h_{rad}

net energy absorbed:

$$q_{net} = A \epsilon \Big|_{T_s} \sigma (T_s^4 - T_{body}^4)$$

assuming: $\epsilon|_{T_s} \approx \epsilon|_{T_b}$

equating with expression for h : $h_{rad} A (T_s - T_b) = A \epsilon \Big|_{T_s} \sigma (T_s^4 - T_{body}^4)$

$$h_{rad} = \frac{\epsilon|_{T_s} \sigma (T_s^4 - T_{body}^4)}{(T_s - T_{body})}$$

Geankoplis 4th ed., eqn 4.10-10 p304

28

© Faith A. Morrison, Michigan Tech U.

Example: Geankoplis 4.10-3

A horizontal oxidized steel pipe carrying steam and having an OD of 0.1683m has a surface temperature of 374.9 K and is exposed to air at 297.1 K in a large enclosure. Calculate the heat loss for 0.305 m of pipe.

$$\varepsilon_{steel} = 0.79$$

29

© Faith A. Morrison, Michigan Tech U.

Example: Geankoplis 4.10-3

A horizontal oxidized steel pipe carrying steam and having an OD of 0.1683m has a surface temperature of 374.9 K and is exposed to air at 297.1 K in a large enclosure. Calculate the heat loss for 0.305 m of pipe.

Answers:

$$h_{radiation} = 6.9W/m^2K$$

$$h_{convection} = 6.1W/m^2K$$

$$Q = 163W$$

$$\varepsilon_{steel} = 0.79$$

30

© Faith A. Morrison, Michigan Tech U.

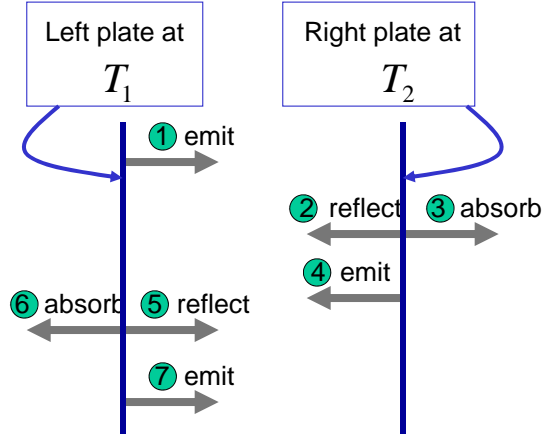
One final topic: **Radiation Heat Transfer Between Two Infinite Plates**

Consider a quantity of radiation energy that is emitted from surface 1.

Reflected= not absorbed

See: Geankoplis, section 4.11B

Also: Bird, Stewart, and Lightfoot, "Transport Phenomena" 1960 Wiley PP446-448



31

© Faith A. Morrison, Michigan Tech U.

First round – surface 2

Radiation Heat Transfer Between Two Infinite Plates

Quantity of energy **incident** at surface 2:

$$\frac{q_{1-2}}{A} = \epsilon_1 \sigma T_1^4$$

Quantity of energy **absorbed** at surface 2:

$$\alpha_2 \left(\frac{q_{1-2}}{A} \right) A = \epsilon_2 (\epsilon_1 \sigma T_1^4) A$$

$\alpha_2 = \epsilon_2$

Quantity of energy **reflected** from surface 2:

$$(1 - \epsilon_2) (\epsilon_1 A \sigma T_1^4)$$

fraction reflected (not absorbed) incident energy

This energy goes back to surface 1.

32

© Faith A. Morrison, Michigan Tech U.

Second round – surface 1

**Radiation Heat Transfer
Between Two Infinite Plates**

Quantity of energy **absorbed** at surface 1
(second round):

Quantity of energy **reflected** from surface 1
(second round):

$$\underbrace{\epsilon_1}_{\text{fraction absorbed}} \underbrace{\left[(1 - \epsilon_2) (\epsilon_1 A \sigma T_1^4) \right]}_{\text{incident energy}}$$

$$\underbrace{(1 - \epsilon_1)}_{\text{fraction reflected (not absorbed)}} \underbrace{\left[(1 - \epsilon_2) (\epsilon_1 A \sigma T_1^4) \right]}_{\text{incident energy}}$$

33
© Faith A. Morrison, Michigan Tech U.

Third round – surface 2

**Radiation Heat Transfer
Between Two Infinite Plates**

Quantity of energy **absorbed** at surface 2
(third round):

Quantity of energy **reflected** from surface 2
(third round):

$$\underbrace{\epsilon_2}_{\text{fraction absorbed}} \underbrace{\left[(1 - \epsilon_1)(1 - \epsilon_2) (\epsilon_1 A \sigma T_1^4) \right]}_{\text{incident energy}}$$

$$\underbrace{(1 - \epsilon_2)}_{\text{fraction reflected (not absorbed)}} \underbrace{\left[(1 - \epsilon_1)(1 - \epsilon_2) (\epsilon_1 A \sigma T_1^4) \right]}_{\text{incident energy}}$$

There is a pattern.

34
© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer Between Two Infinite Plates

Now, calculate the radiation energy going from surface 1 to surface 2:

Later, calculate energy from 2 to 1; then subtract to obtain **net energy transferred**.

$$\begin{aligned}
 q_{1-2} &= \left(\begin{array}{c} \text{energy from} \\ 1 \rightarrow 2 \end{array} \right) = \sum \left(\begin{array}{c} \text{energy absorbed} \\ \text{at surface 2} \end{array} \right) \\
 &= \varepsilon_2 (\varepsilon_1 A \sigma T_1^4) \\
 &\quad + \varepsilon_2 (1 - \varepsilon_1)(1 - \varepsilon_2) (\varepsilon_1 A \sigma T_1^4) \\
 &\quad + \varepsilon_2 (1 - \varepsilon_1)^2 (1 - \varepsilon_2)^2 (\varepsilon_1 A \sigma T_1^4) \\
 &\quad \dots + \varepsilon_2 (1 - \varepsilon_1)^n (1 - \varepsilon_2)^n (\varepsilon_1 A \sigma T_1^4) + \dots
 \end{aligned}$$

35

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer Between Two Infinite Plates

Radiation energy going from surface 1 to surface 2:

$$q_{1-2} = \varepsilon_1 \varepsilon_2 A \sigma T_1^4 \sum_{n=0}^{\infty} (1 - \varepsilon_1)^n (1 - \varepsilon_2)^n$$

How can we calculate $\sum_{n=0}^{\infty} x^n$?

Answer: $1/(1-x)$

36

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer
Between Two Infinite Plates

Radiation energy going from
surface 1 to surface 2:

$$q_{1-2} = \frac{\epsilon_1 \epsilon_2 A \sigma T_1^4}{1 - [(1 - \epsilon_1)(1 - \epsilon_2)]}$$

$$= \frac{\epsilon_1 \epsilon_2 A \sigma T_1^4}{1 - [1 - \epsilon_1 - \epsilon_2 + \epsilon_1 \epsilon_2]} = \frac{\epsilon_1 \epsilon_2 A \sigma T_1^4}{\epsilon_1 + \epsilon_2 - \epsilon_1 \epsilon_2}$$

$$\frac{q_{1-2}}{A} = \frac{\sigma T_1^4}{\frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} - 1}$$

Final Result

37

© Faith A. Morrison, Michigan Tech U.

Radiation Heat Transfer
Between Two Infinite Plates

Radiation energy going from
surface 1 to surface 2:

$$\frac{q_{1-2}}{A} = \frac{\sigma T_1^4}{\frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} - 1}$$

Radiation energy going from
surface 2 to surface 1:

$$\frac{q_{2-1}}{A} = \frac{\sigma T_2^4}{\frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} - 1}$$

NET Radiation energy going
from surface 1 to surface 2:

$$\frac{q_{1-2} - q_{2-1}}{A} = \frac{\sigma (T_1^4 - T_2^4)}{\left(\frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} - 1 \right)}$$

38

© Faith A. Morrison, Michigan

Radiation Shields

Purpose of Heat Shields:
To reduce the amount of energy transfer from (hotter) plate at T_1 to second (cooler) plate at T_3 .

Note:
 $q_{net,1 \rightarrow 2} = q_{net,2 \rightarrow 3} = q$

39
© Faith A. Morrison, Michigan Tech U.

Analysis of Radiation Shields

We will assume that the emissivity is the same for all surfaces.

$$\frac{q_{net,1 \rightarrow 2}}{A} = \frac{\sigma(T_1^4 - T_2^4)}{\left(\frac{1}{\epsilon} + \frac{1}{\epsilon} - 1\right)}$$

$$\frac{q_{net,2 \rightarrow 3}}{A} = \frac{\sigma(T_2^4 - T_3^4)}{\left(\frac{1}{\epsilon} + \frac{1}{\epsilon} - 1\right)}$$

Now we eliminate T_2 between these equations.

Note:
 $q_{net,1 \rightarrow 2} = q_{net,2 \rightarrow 3} = q$

40
© Faith A. Morrison, Michigan Tech U.

Analysis of Radiation Shields

$$\frac{q}{A} = \frac{\sigma(T_1^4 - T_2^4)}{\left(\frac{2}{\epsilon} - 1\right)}$$

$$\frac{q}{A} = \frac{\sigma(T_2^4 - T_3^4)}{\left(\frac{2}{\epsilon} - 1\right)}$$

$$T_2^4 = \frac{q}{\sigma A} \left(\frac{2}{\epsilon} - 1\right) + T_3^4$$

$$\frac{q}{\sigma A} \left(\frac{2}{\epsilon} - 1\right) = T_1^4 - \frac{q}{\sigma A} \left(\frac{2}{\epsilon} - 1\right) - T_3^4$$

$$\frac{2q}{\sigma A} \left(\frac{2}{\epsilon} - 1\right) = T_1^4 - T_3^4$$

$$\frac{q}{A} = \left(\frac{1}{2}\right) \frac{\sigma(T_1^4 - T_3^4)}{(2/\epsilon - 1)}$$

Radiation Shield

T_1 T_2 T_3

41
© Faith A. Morrison, Michigan Tech U.

Analysis of Radiation Shields

1 Heat Shield

$$\frac{q}{A} = \left(\frac{1}{2}\right) \frac{\sigma(T_1^4 - T_3^4)}{\left(\frac{2}{\epsilon} - 1\right)}$$

With one heat shield present, q falls by half compared to no heat shield.

by the same analysis,

N Heat Shields

$$\frac{q}{A} = \left(\frac{1}{N+1}\right) \frac{\sigma(T_1^4 - T_3^4)}{\left(\frac{2}{\epsilon} - 1\right)}$$

With N heat shields present, q falls by a factor of $1/N$ compared to no heat shield.

Radiation Shield

T_1 T_2 T_3

42
© Faith A. Morrison, Michigan Tech U.

Radiation

Summary:

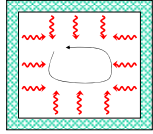
$$\sigma = 0.1712 \times 10^{-8} \frac{BTU}{h \text{ ft}^2 R^4}$$

$$\sigma = 5.676 \times 10^{-8} \frac{W}{m^2 K^4}$$

General properties:

- Absorptivity, α
 - gray body: $\alpha = \text{constant}$
 - black body: $\alpha = 1$
- Emissivity, ε
 $q_{emit} = \varepsilon q_{emit, blackbody}$
- Kirchoff's law: $\alpha = \varepsilon$
- Stefan-Boltzman law
 $\frac{q_{emit, blackbody}}{A} = \sigma T^4$

Heat transfer coefficient:




$$h_{rad} = \frac{\varepsilon |T_s \sigma (T_s^4 - T_{body}^4)}{(T_s - T_{body})}$$

Geankoplis 4th ed., eqn 4.10-10 p304

Heat shields:

$$\frac{q}{A} = \left(\frac{1}{N+1} \right) \frac{\sigma (T_1^4 - T_3^4)}{\left(\frac{2}{\varepsilon} - 1 \right)}$$




Always use absolute temperature (Kelvin) in radiation calculations.

43

© Faith A. Morrison, Michigan Tech U.

CM3110

Transport Processes and Unit Operations I




Michigan Tech

Part 2:

Professor Faith Morrison

Department of Chemical Engineering
Michigan Technological University

CM3110 - Momentum and Heat Transport
CM3120 - Heat and Mass Transport



www.chem.mtu.edu/~fmorriso/cm310/cm310.html

44

© Faith A. Morrison, Michigan Tech U.

CM3110
Transport Processes and Unit Operations I
Part 2: Heat Transfer

Summary

Within homogeneous phases:

- Microscopic Energy Balances
- 1D Steady solutions

rectangular:

$$\frac{q_x}{A} = C_1$$

$$T = ax + b$$

cylindrical:

$$\frac{q_r}{A} = \frac{C_1}{r}$$

$$T = a \ln x + b$$

$$\rho \hat{c}_p \left(\frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S_e$$

- Temperature and *Newton's law of cooling* boundary conditions (if h is supplied)
- Unsteady solutions (from literature)
 - ✓ Carslaw and Jeager
 - ✓ Heisler charts

45

© Faith A. Morrison, Michigan Tech U.

CM3110
Transport Processes and Unit Operations I
Part 2: Heat Transfer

Summary

Across phase boundaries:

- Microscopic Energy, Momentum, and Mass Balances

Micro momentum:

$$\rho \left(\frac{\partial \underline{v}}{\partial t} + \underline{v} \cdot \nabla \underline{v} \right) = -\nabla p + \mu \nabla^2 \underline{v} + \rho \underline{g}$$

Micro energy:

$$\rho \hat{c}_p \left(\frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \right) = k \nabla^2 T + S_e$$

- Simultaneous effects (complex)
- Solutions are difficult to obtain (and often not really necessary)
 - use *dimensional analysis* to obtain h
- h Data correlations for:
 - ✓ forced convection,
 - ✓ natural convection
 - ✓ evaporation/condensation
 - ✓ radiation

} (use in design)

46

© Faith A. Morrison, Michigan Tech U.

CM3110
Transport Processes and Unit Operations I
Part 2: Heat Transfer



Michigan Tech

Summary

Heat Transfer Unit Operations

- Macroscopic energy balances
- Heat Exchangers
 - ✓ double pipe (ΔT_{lm})
 - ✓ Shell-and-tube
 - ✓ Heat exchanger effectiveness
- Evaporators/ Condensers
- Ovens
- Heat Shields

47

© Faith A. Morrison, Michigan Tech U.

CM3110
Transport Processes and Unit Operations I

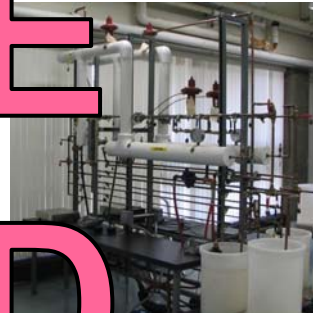


Michigan Tech

Professor Faith Morrison

Department of Chemical Engineering
Michigan Technological University

CM3110 - Momentum and Heat Transport
CM3120 - Heat and Mass Transport



**THE
END**

www.chem.mtu.edu/~fmorriso/cm310/cm310.html

48

© Faith A. Morrison, Michigan Tech U.