

Final Exam
SOLN
CM 4650
2018

①

1. $\underline{a} \cdot \nabla \underline{b}$

$$a_i \hat{e}_i \cdot \frac{\partial}{\partial x_p} \hat{e}_p b_m \hat{e}_m$$

δ_{ip} "i becomes p"

$$= a_p \frac{\partial b_m}{\partial x_p} \hat{e}_m \quad \underline{\text{vector}}$$

$$= \sum_{p=1}^3 \sum_{m=1}^3 a_p \frac{\partial b_m}{\partial x_p} \hat{e}_m$$

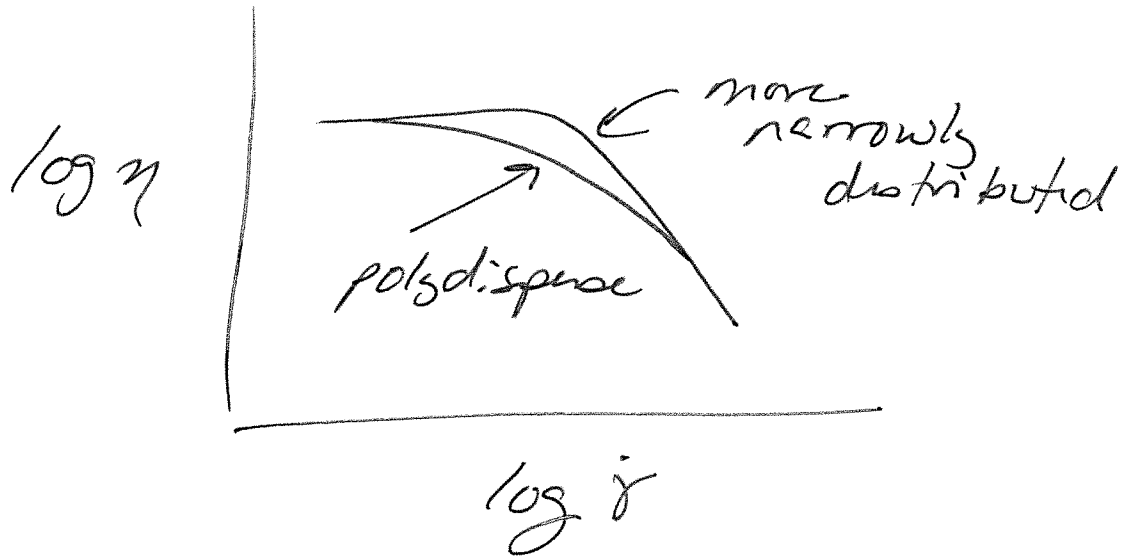
$$= \begin{pmatrix} \sum_{p=1}^3 a_p \frac{\partial b_1}{\partial x_p} \\ \sum_{p=1}^3 a_p \frac{\partial b_2}{\partial x_p} \\ \sum_{p=1}^3 a_p \frac{\partial b_3}{\partial x_p} \end{pmatrix}_{123} //$$

2.

- a) False. Cox Merz is only true when it's true, i.e. it must be shown to hold before we use it.
- b) False. Troutman's rule generally does not hold at large γ .
- c) We can tell if a material is entangled
- ① from the "hockey stick graph" i.e. $\log \eta$ vs $\log M$, which shows a break at M_c , the critical MW for entanglement
 - ② In SAS $G'(\omega)$ has a plateau $\sim G' \sim 10^4 - 10^5$ when entangled.

(3)

d)



3. What is ψ_1 for Lodge model? ④

Lodge

$$\underline{\underline{\tau}} = - \int_{-\infty}^t \frac{\eta_0}{\lambda^2} e^{\frac{-(t-t')}{\lambda}} \underline{\underline{C}}^{-1}(t', t) dt'$$

For χ_1 , $\dot{\gamma}(t'') = \dot{\gamma}_0$

$$\psi_1 = \frac{-(\tau_{11} - \tau_{22})}{\dot{\gamma}_0^2}$$

Table 9.3

$$\underline{\underline{C}}' = \begin{pmatrix} 1 + \gamma^2 & \gamma & 0 \\ \gamma & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

We need γ ; $\gamma = \int_{t'}^t \dot{\gamma}(t'') dt''$

$$= \int_{t'}^t \dot{\gamma}_0 dt''$$

$$= \dot{\gamma}_0 t'' \Big|_{t'}^t = \boxed{\dot{\gamma}_0 (t - t')} = \gamma$$

(5)

We need $\tau_{11} - \tau_{22}$

$$-\tau_{11} = \int_{-\infty}^t \frac{\gamma_0}{\lambda^2} e^{\frac{-(t-t')}{\lambda}} (1+\delta^2) dt'$$

$$-\tau_{22} = \int_{-\infty}^t \frac{\gamma_0}{\lambda^2} e^{\frac{-(t-t')}{\lambda}} (1) dt'$$

$$-(\tau_{11} - \tau_{22}) = \int_{-\infty}^t \frac{\gamma_0}{\lambda^2} e^{\frac{-(t-t')}{\lambda}} \underbrace{(1+\delta^2)}_{\delta^2 = \dot{\gamma}_0^2 (t-t')^2} dt'$$

$$\frac{-(\tau_{11} - \tau_{22}) \lambda^2}{\gamma_0 \dot{\gamma}_0^2} = \lambda^2 \lambda \int_{-\infty}^t \underbrace{e^{\frac{-(t-t')}{\lambda}}}_{= e^u} \underbrace{\frac{(t-t')^2}{\lambda^2}}_{u^2} \underbrace{\frac{dt'}{\lambda}}_{du}$$

$$\Rightarrow u = \frac{-t}{\lambda} + \frac{t'}{\lambda} = -\frac{1}{\lambda}(t-t')$$

$$\frac{du}{dt'} = \frac{1}{\lambda}$$

$$\Rightarrow \frac{(t-t')^2}{\lambda^2} = u^2$$

$$\rightarrow du = \frac{1}{\lambda} dt'$$

$$\begin{array}{ll} t' = -\infty & u = -\infty \\ t' = t & u = 0 \end{array}$$

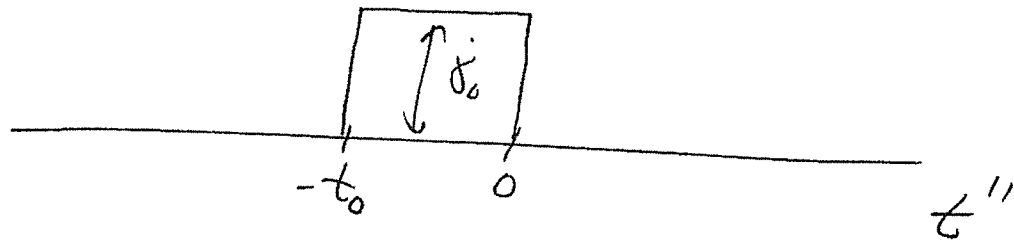
(6)

$$\begin{aligned}
 \frac{-(\tau_{11} - \tau_{22})}{\eta_0 \lambda \dot{\gamma}_0^2} &= \int_{-\infty}^0 e^u u^2 du \\
 &= e^u (u^2 - 2u + 2) \Big|_{-\infty}^0 \\
 &= 2 - 0 = 2
 \end{aligned}$$

$$\psi_1 = \frac{-(\tau_{11} - \tau_{22})}{\dot{\gamma}_0^2} = \boxed{2\eta_0\lambda}$$

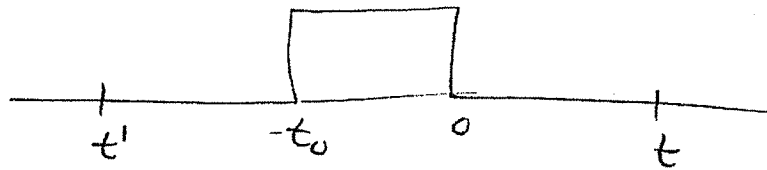
(7)

7.



$$\gamma_{21}(t', t) = \int_{t'}^t \dot{J}(t'') dt''$$

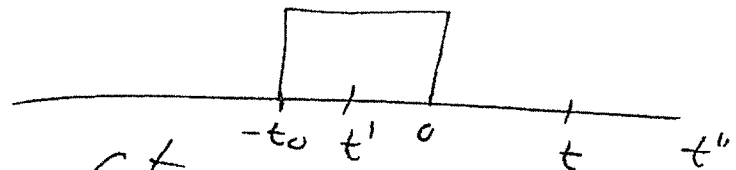
CASE I: $t' < -t_0$



$$\gamma = \int_{t'}^{-t_0} 0 dt'' + \int_{-t_0}^0 j_0 dt'' + \int_0^t 0 dt''$$

$$\boxed{\gamma = t_0 j_0}$$

CASE II: $-t_0 \leq t' \leq 0$



$$\gamma = \int_{t'}^0 j_0 dt'' + \int_0^t 0 dt''$$

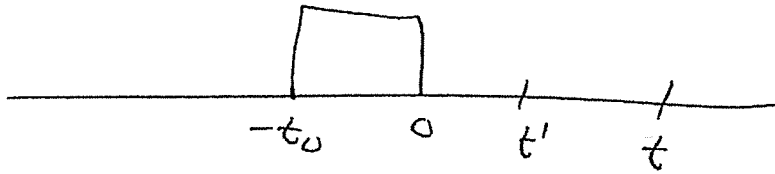
$$\boxed{\gamma = -j_0 t'}$$

(Note $t' < 0 \therefore \gamma$ positive)

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CASE III

$$t' > 0$$

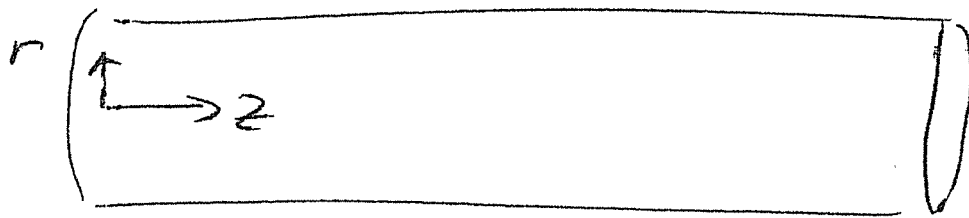


$$\gamma = \int_{t'}^t 0 dt'' = 0$$

$$\gamma = \begin{cases} t_0 \ddot{\gamma}_0 & t' < -t_0 \\ -\ddot{\gamma}_0 t' & -t_0 \leq t' \leq 0 \\ 0 & t' > 0 \end{cases}$$

//

5. PL-GNF



$$\rho \left(\underbrace{\frac{\partial \underline{v}}{\partial t}}_{\text{steady}} + \underbrace{\underline{v} \cdot \nabla \underline{v}}_{\text{uni dir}} \right) = -\nabla p - \nabla \cdot \underline{\underline{\tau}} + \underbrace{\rho g}_{\text{neglect}}$$

$$\underline{v} = \begin{pmatrix} v_r \\ v_\theta \\ v_z \end{pmatrix}_{r\theta z} = \begin{pmatrix} 0 \\ 0 \\ v_z \end{pmatrix}$$

$$\nabla \cdot \underline{v} = 0 \Rightarrow \boxed{\frac{\partial v_z}{\partial z} = 0} \quad \begin{matrix} \text{see} \\ \text{next} \\ \text{pg} \end{matrix}$$

$$\underline{\underline{\tau}} = \eta(\dot{\gamma}) \dot{\underline{\underline{\gamma}}} = \eta(\dot{\gamma}) \left(\nabla \underline{v} + (\nabla \underline{v})^T \right)$$

↑
(see next pg)

$\eta = m \dot{\gamma}^{n-1}$

C.2 Differential Operations in Curvilinear Coordinates

TABLE C.3
Differential Operations in the Cylindrical Coordinate System r, θ, z

$$\underline{w} = \begin{pmatrix} w_r \\ w_\theta \\ w_z \end{pmatrix}_{r\theta z}$$

$$\nabla = \hat{e}_r \frac{\partial}{\partial r} + \hat{e}_\theta \frac{1}{r} \frac{\partial}{\partial \theta} + \hat{e}_z \frac{\partial}{\partial z}$$

$$\nabla a = \begin{pmatrix} \frac{\partial a}{\partial r} \\ \frac{1}{r} \frac{\partial a}{\partial \theta} \\ \frac{\partial a}{\partial z} \end{pmatrix}_{r\theta z}$$

$$\nabla \cdot \nabla a = \nabla^2 a = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial a}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 a}{\partial \theta^2} + \frac{\partial^2 a}{\partial z^2}$$

$$\nabla \cdot \underline{w} = \frac{1}{r} \frac{\partial}{\partial r} (r w_r) + \frac{1}{r} \frac{\partial w_\theta}{\partial \theta} + \frac{\partial w_z}{\partial z} = 0$$

$$\nabla \times \underline{w} = \begin{pmatrix} \frac{1}{r} \frac{\partial w_z}{\partial \theta} - \frac{\partial w_\theta}{\partial z} \\ \frac{\partial w_r}{\partial z} - \frac{\partial w_z}{\partial r} \\ \frac{1}{r} \frac{\partial (r w_\theta)}{\partial r} - \frac{1}{r} \frac{\partial w_r}{\partial \theta} \end{pmatrix}_{r\theta z}$$

$$\Rightarrow \frac{\partial w_z}{\partial z} = 0$$

$$\underline{A} = \begin{pmatrix} A_{rr} & A_{r\theta} & A_{rz} \\ A_{\theta r} & A_{\theta\theta} & A_{\theta z} \\ A_{zr} & A_{z\theta} & A_{zz} \end{pmatrix}_{r\theta z}$$

$$\nabla \underline{w} = \begin{pmatrix} \frac{\partial w_r}{\partial r} & \frac{\partial w_\theta}{\partial r} & \frac{\partial w_z}{\partial r} \\ \frac{1}{r} \frac{\partial w_r}{\partial \theta} - \frac{w_\theta}{r} & \frac{1}{r} \frac{\partial w_\theta}{\partial \theta} + \frac{w_r}{r} & \frac{1}{r} \frac{\partial w_z}{\partial \theta} \\ \frac{\partial w_r}{\partial z} & \frac{\partial w_\theta}{\partial z} & \frac{\partial w_z}{\partial z} \end{pmatrix}_{r\theta z}$$

$$\nabla^2 \underline{w} = \begin{pmatrix} \frac{\partial}{\partial r} \left[\frac{1}{r} \frac{\partial (r w_r)}{\partial r} \right] + \frac{1}{r^2} \frac{\partial^2 w_r}{\partial \theta^2} + \frac{\partial^2 w_r}{\partial z^2} - \frac{2}{r^2} \frac{\partial w_\theta}{\partial \theta} \\ \frac{\partial}{\partial r} \left[\frac{1}{r} \frac{\partial (r w_\theta)}{\partial r} \right] + \frac{1}{r^2} \frac{\partial^2 w_\theta}{\partial \theta^2} + \frac{\partial^2 w_\theta}{\partial z^2} + \frac{2}{r^2} \frac{\partial w_r}{\partial \theta} \\ \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial w_z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 w_z}{\partial \theta^2} + \frac{\partial^2 w_z}{\partial z^2} \end{pmatrix}_{r\theta z}$$

for $\underline{w} = \underline{v}$

for $\underline{w} = \underline{v}$

Table C.3

(11)

$$\underline{\nabla} V = \begin{pmatrix} \cancel{\frac{\partial V_r}{\partial r}} & \cancel{\frac{\partial V_\theta}{\partial r}} & \frac{\partial V_z}{\partial r} \\ \cancel{\frac{1}{r} \frac{\partial V_r}{\partial \theta} - \frac{V_\theta}{r}} & \cancel{\frac{1}{r} \frac{\partial V_\theta}{\partial \theta} + \frac{V_r}{r}} & \cancel{\frac{1}{r} \frac{\partial V_z}{\partial \theta}} \\ \cancel{\frac{\partial V_r}{\partial z}} & \cancel{\frac{\partial V_\theta}{\partial z}} & \cancel{\frac{\partial V_z}{\partial z}} \end{pmatrix} \begin{matrix} \oplus \\ s_z \hat{e}_z \\ r \hat{e}_z \end{matrix}$$

$\nabla \cdot \underline{V} = 0$

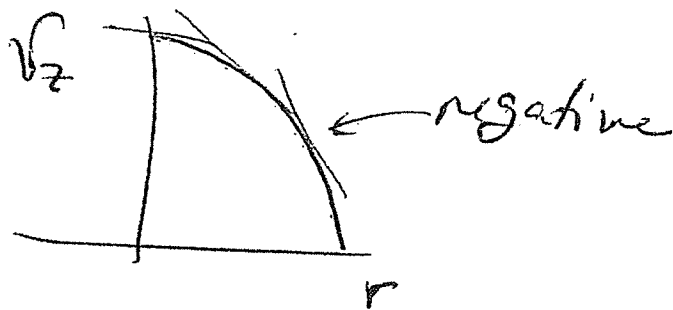
$$= \begin{pmatrix} 0 & 0 & \frac{\partial V_z}{\partial r} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} r \hat{e}_z$$

$$\underline{\dot{\gamma}} = \begin{pmatrix} 0 & 0 & \frac{\partial V_z}{\partial r} \\ 0 & 0 & 0 \\ \frac{\partial V_z}{\partial r} & 0 & 0 \end{pmatrix} r \hat{e}_z$$

(12)

$$|\dot{\underline{x}}| = \sqrt{\frac{\dot{\underline{x}} \cdot \dot{\underline{x}}}{2}} = \sqrt{\left(\frac{\partial V_z}{\partial r}\right)^2}$$

$$= \left| \frac{\partial V_z}{\partial r} \right| = \left(\frac{\partial V_z}{\partial r} \right)$$



$$\underline{\tau} = -\eta \dot{\underline{x}} = -m \dot{\underline{x}}^{n-1} \dot{\underline{x}}$$

$$= -m \left(-\frac{\partial V_z}{\partial r} \right)^{n-1} \begin{pmatrix} 0 & 0 & \frac{\partial V_z}{\partial r} \\ 0 & 0 & 0 \\ \frac{\partial V_z}{\partial r} & 0 & 0 \end{pmatrix}_{rot}$$

$$= \begin{pmatrix} \tau_{rr} & \tau_{r\theta} & \tau_{rz} \\ \tau_{\theta r} & \tau_{\theta\theta} & \tau_{\theta z} \\ \tau_{zr} & \tau_{z\theta} & \tau_{zz} \end{pmatrix}_{rot}$$

Back to Cauchy Momentum Eqn

(B)

$$0 = -\nabla p - \underbrace{\nabla \cdot \underline{\underline{\tau}}}_{\text{and from Table C.3}}$$

(see next page)

$$\begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}_{r\theta z} = \begin{pmatrix} -\frac{\partial p}{\partial r} \\ -\frac{1}{r} \frac{\partial p}{\partial \theta} \\ -\frac{\partial p}{\partial z} \end{pmatrix}_{r\theta z} - \begin{pmatrix} \frac{\partial}{\partial z} (\cancel{\underline{\tau}_{rz}}) \\ 0 \\ \frac{1}{r} \frac{\partial}{\partial r} (r \underline{\tau}_{rz}) \end{pmatrix}_{r\theta z}$$

v_z does not vary w/ z

Now we consider each component individually $\rightarrow$

for
 $\underline{A} = \underline{\underline{A}}$

$$\nabla \cdot \underline{A} = \begin{pmatrix} \frac{1}{r} \frac{\partial}{\partial r} (r A_{rr}) + \frac{1}{r} \frac{\partial A_{\theta r}}{\partial \theta} + \frac{\partial A_{zr}}{\partial z} - \frac{A_{\theta\theta}}{r} \\ \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_{r\theta}) + \frac{1}{r} \frac{\partial A_{\theta\theta}}{\partial \theta} + \frac{\partial A_{z\theta}}{\partial z} + \frac{A_{\theta r} - A_{r\theta}}{r} \\ \frac{1}{r} \frac{\partial}{\partial r} (r A_{rz}) + \frac{1}{r} \frac{\partial A_{\theta z}}{\partial \theta} + \frac{\partial A_{zz}}{\partial z} \end{pmatrix}_{r\theta z} \quad (C.3-10)$$

$$\underline{u} \cdot \nabla \underline{w} = \begin{pmatrix} u_r \left(\frac{\partial w_r}{\partial r} \right) + u_\theta \left(\frac{1}{r} \frac{\partial w_r}{\partial \theta} - \frac{w_\theta}{r} \right) + u_z \left(\frac{\partial w_r}{\partial z} \right) \\ u_r \left(\frac{\partial w_\theta}{\partial r} \right) + u_\theta \left(\frac{1}{r} \frac{\partial w_\theta}{\partial \theta} + \frac{w_r}{r} \right) + u_z \left(\frac{\partial w_\theta}{\partial z} \right) \\ u_r \left(\frac{\partial w_z}{\partial r} \right) + u_\theta \left(\frac{1}{r} \frac{\partial w_z}{\partial \theta} \right) + u_z \left(\frac{\partial w_z}{\partial z} \right) \end{pmatrix}_{r\theta z} \quad (C.3-11)$$

TABLE C.4

Differential Operations in the Spherical Coordinate System r, θ, ϕ

$$\underline{w} = \begin{pmatrix} w_r \\ w_\theta \\ w_\phi \end{pmatrix}_{r\theta\phi} \quad (C.4-1)$$

$$\nabla = \hat{e}_r \frac{\partial}{\partial r} + \hat{e}_\theta \frac{1}{r} \frac{\partial}{\partial \theta} + \hat{e}_\phi \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \quad (C.4-2)$$

$$\nabla a = \begin{pmatrix} \frac{\partial a}{\partial r} \\ \frac{1}{r} \frac{\partial a}{\partial \theta} \\ \frac{1}{r \sin \theta} \frac{\partial a}{\partial \phi} \end{pmatrix}_{r\theta\phi} \quad (C.4-3)$$

$$\nabla \cdot \nabla a = \nabla^2 a = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial a}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial a}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 a}{\partial \phi^2} \quad (C.4-4)$$

$$\nabla \cdot \underline{w} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 w_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (w_\theta \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial w_\phi}{\partial \phi} \quad (C.4-5)$$

$$\nabla \times \underline{w} = \begin{pmatrix} \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (w_\phi \sin \theta) - \frac{1}{r \sin \theta} \frac{\partial w_\theta}{\partial \phi} \\ \frac{1}{r \sin \theta} \frac{\partial w_r}{\partial \phi} - \frac{1}{r} \frac{\partial}{\partial r} (r w_\phi) \\ \frac{1}{r} \frac{\partial}{\partial r} (r w_\theta) - \frac{1}{r} \frac{\partial w_r}{\partial \theta} \end{pmatrix}_{r\theta\phi} \quad (C.4-6)$$

$$\underline{A} = \begin{pmatrix} A_{rr} & A_{r\theta} & A_{r\phi} \\ A_{\theta r} & A_{\theta\theta} & A_{\theta\phi} \\ A_{\phi r} & A_{\phi\theta} & A_{\phi\phi} \end{pmatrix}_{r\theta\phi} \quad (C.4-7)$$

continued

r -component:

$$\frac{\partial P}{\partial r} = 0$$

θ -component:

$$\frac{1}{r} \frac{\partial P}{\partial \theta} = 0 \quad \Rightarrow \quad \frac{\partial P}{\partial \theta} = 0$$

z -component:

$$\lambda = \underbrace{\frac{\partial P}{\partial z}}_{f(z)} = - \underbrace{\frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rz})}_{f(r)}$$

$$\frac{\partial P}{\partial z} = \lambda$$

$$P = \lambda z + C_1$$

$$\text{BC: } \begin{aligned} z=0 & \quad P=P_0 \\ z=L & \quad P=P_L \end{aligned}$$

$$C_1 = P_0$$

$$\lambda = \frac{P_L - P_0}{L}$$

$$P = \frac{P_L - P_0}{L} z + P_0$$

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$$-\frac{\Delta P}{L} = -\frac{1}{r} \frac{d}{dr} (r \tau_{rz})$$

$$\frac{\Delta P r}{L} = \frac{d}{dr} (r \tau_{rz})$$

$$r \tau_{rz} = \frac{\Delta P}{L} \frac{r^2}{2} + C_2$$

$$\tau_{rz} = \frac{\Delta P}{2L} r + \frac{C_2}{r}$$

$$\text{BL: } r=0$$

$$\tau_{rz} = \text{finite}$$

$$\Rightarrow C_2 = 0$$

$$\tau_{rz} = \frac{\Delta P r}{2L}$$

From constitutive eqn:

$$-m \left(-\frac{\partial V_z}{\partial r} \right)^{n-1} \left(\frac{\partial V_z}{\partial r} \right) = \frac{\Delta P r}{2L}$$

$$\left(-\frac{\partial V_z}{\partial r} \right)^n = \frac{\Delta P r}{2Lm}$$

$$-\frac{dV_z}{dr} = \left(\frac{\Delta P}{2mL} \right)^{\frac{1}{n}} r^{\frac{1}{n}}$$

$$V_z = - \left(\frac{\Delta P}{2mL} \right)^{\frac{1}{n}} \frac{r^{\frac{1}{n}+1}}{\frac{1}{n}+1} + C_3$$

BC: $r=R$ $V_z=0$

(already used $r=0$
BC but could

say $\frac{dV_z}{dr}=0$ there //